

Technical Memorandum 26

# A simple computer model for \_\_\_\_\_\_ terrestrial and solar radiation transfer

I.M. Vardavas and L.M. Cannon

Supervising Scientist for the Alligator Rivers Region

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**TECHNICAL MEMORANDUM 26** 

# A SIMPLE COMPUTER MODEL FOR TERRESTRIAL AND SOLAR RADIATION TRANSFER

Ilias M. Vardavas and Lisa M. Cannon

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# ABSTRACT

Ilias M. Vardavas & Lisa M. Cannon (1989). A simple computer model for terrestrial and solar radiation transfer. Technical Memorandum 26, Supervising Scientist for the Alligator Rivers Region.

A simple radiative-convective atmospheric model is presented for rapidly computing the solar and terrestrial fluxes at the top of the atmosphere and at the ground. The model parameters are measurable meteorological quantities with water vapour playing a key role in the determination of the ratio of infrared flux emitted by a water surface to the absorbed solar flux; this infrared fraction is important in Heat Balance evaporation models. The model and its computer programme can be used to examine, for example, the response of the global mean surface temperature to changes in  $CO_2$ , cloud cover and solar-constant.

# EXPLANATORY NOTES

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# Pressure Levels

p<sub>g</sub> = ground pressure p<sub>c</sub> = cloud top pressure p<sub>cb</sub> = cloud base pressure

# Atmosphere IR Optical Depths and Transmissivities

$\tau_{\mathbf{g}}$	=	$\tau(0,p_g) = \tau(0,p_g) = -$	total optical depth above cloud optical depth
$\tau_{cb}$	=	$\tau(\mathbf{p}_{cb},\mathbf{p}_{g}) =$	below cloud optical depth
t <sub>g</sub> t <sub>c</sub> t <sub>cb</sub>	=	$\exp(-\tau_{\rm g}) = \\ \exp(-\tau_{\rm c}) = \\ \exp(-\tau_{\rm cb}) =$	total transmissivity above cloud transmissivity below cloud transmissivity
t <sub>1</sub> t <sub>2</sub> t <sub>3</sub>	H	$exp(-\tau(p,0))exp(-\tau(p,p_g))exp(-\tau(p,p_{cb}))$	<ul> <li>transmissivity between pressure level p and top of atmosphere</li> <li>transmissivity between pressure level p and ground</li> <li>transmissivity between pressure level p and cloud base (p&gt;p<sub>cb</sub>)</li> </ul>
$t_4$	=	$exp(-\tau(p,p_c))$	= transmissivity between p and cloud top $(p < p_c)$

# **1 INTRODUCTION**

Recently, Vardavas & Carver (1984*a,b*) developed a one dimensional (1-D) radiativeconvective modelling technique for rapidly computing the mean vertical atmospheric temperature profile. They presented parameterizations for both solar (~ 0.1-5.0  $\mu$ m) and terrestrial (greater than about 5  $\mu$ m) radiation transfer. The infrared, visible and ultraviolet regions were divided into fine spectral intervals within which atmospheric molecules of water vapour (H<sub>2</sub>O), carbon dioxide (CO<sub>2</sub>), ozone (O<sub>3</sub>), methane (CH<sub>4</sub>) and ammonia (NH<sub>3</sub>) have important absorption bands.

This fine spectral radiation transfer was devised in order to model climatic change due to spectral variations in the solar flux and changes in the terrestrial surface albedo (or reflectivity). The model has been used to examine the response of the earth's global mean vertical temperature to changes in atmospheric composition (e.g. changes in  $CO_2$  and  $CH_4$ ), cloud cover and solar-constant (Vardavas & Carver 1985), and to solar cycle uv flux variations (Vardavas 1987*a*).

In the present work, a simpler version (SV) of the above more complex model (CV) is presented. It is based on simple expressions for the total absorption of infrared, visible and ultraviolet radiation by the atmospheric molecules. This allows the atmosphere to be represented as layers having simple transmission coefficients (or transmissivities) which depend on the amount of absorbing molecules in each layer.

The present simple model can be used to compute rapidly solar and terrestrial radiation fluxes at the top of the atmosphere and at the ground using measurable meteorological quantities as model input parameters. It can be easily included in Heat Balance models to compute the net all-wave flux absorbed by a water surface (Vardavas 1987b) while the mean global surface temperature can be readily computed from flux balance between net all-wave incoming and outgoing radiation fluxes at the top of the atmosphere. The present model can thus be used to examine the effects of climatic change due, for example, to changes in atmospheric  $CO_2$ , cloud cover and solar-constant.

The atmospheric model is described in Section 2. The 1-D vertical temperature structure consists of a troposphere with a temperature structure determined by a specified adiabatic lapse rate while the lower stratosphere is taken to be isothermal with a temperature fixed at the tropopause temperature. In Section 3 the evaluation of the terrestrial flux is described both for a clear and cloudy sky. Expressions are derived for computing the net infrared flux at the ground and at the top of the atmosphere. In Section 4 the solar radiation transfer is described and an expression is derived for the planetary albedo which determines the net incoming solar flux at the top of the atmosphere and the net solar flux incident on the ground. Model flux and temperature computations using the present SV model are compared with those of the complex CV model in Section 5. The computer programme and its documentation are given in Section 6. Two sample computes the radiation fluxes at the top of the atmosphere computes the radiation fluxes at the top of the atmosphere computes the radiation fluxes at the top of the atmosphere computer runs are included; one computes the mean global surface temperature, the other computes the radiation fluxes at the top of the atmosphere and at the ground for a specified atmospheric temperature profile.

# **2** ATMOSPHERIC MODEL

#### 2.1 Radiative equilibrium

In the present 1-D radiative-convective model, flux balance at the top of the atmosphere is used to determine the earth's mean global surface temperature,  $T_g$ . The earth is taken to be

in radiative equilibrium and so the net flux at the top of the atmosphere is set equal to zero (see Fig. 1):

$$\overline{F}_{ir}^{\dagger}(\infty) - \overline{F}_{o}^{\downarrow}(\infty) = 0$$
 (1)

where  $\overline{F}_{ir}^{\uparrow}(\infty)$  is the net outgoing terrestrial infrared flux (Wm<sup>-2</sup>) and  $\overline{F}_{o}^{\downarrow}(\infty)$  is the net incoming solar flux at the top of atmosphere. The terrestrial flux is a strong function of the vertical atmospheric temperature structure and the atmospheric infrared opacity. The surface temperature of the earth,  $T_g$ , is obtained iteratively from the flux balance equation. The solar flux absorbed by the atmosphere is weakly dependent on the atmospheric temperature profile but is strongly dependent on the atmospheric gases.

Incoming solar ↓	_ Reflected solar ↑	Outgoing terrestrial †	н	0
	•	·		

Top of atmosphere

#### Terrestrial surface

Figure 1. Flux balance at the top of the atmosphere

#### 2.2 Radiation fields

Incident solar radiation is either absorbed by the surface-atmosphere system or is reflected back to space. The absorbed solar radiation is re-emitted by the surface-atmosphere as terrestrial radiation.

The atmospheric absorption and emission of the solar and terrestrial radiation fields can be treated separately because the flux emitted by the sun consists mainly of short-wave radiation (~ 0.1-5.0  $\mu$ m) while the flux emitted by the earth consists of infrared radiation (> 5  $\mu$ m). Idealized spectral distributions of the two radiation fields are shown in Fig. 2 (Note the difference in units between solar and terrestrial fields). The distributions were generated by Planck functions with T = 5770 K for the sun and T = 252 K for the earth, corresponding to the effective black body temperatures of the two bodies.



Figure 2. The spectral distributions of solar and terrestrial radiation fields

#### 2.3 Greenhouse effect of the atmosphere

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A AN PARTY AND ADDRESS OF A DESCRIPTION

The atmosphere is an envelope of gas and clouds which surrounds the earth and which greatly affects the earth's temperature. Solar radiation incident on the surface is reflected, absorbed and reemitted. Without an atmosphere all of the re-emitted (or terrestrial) radiation would escape to space. However, the presence of an atmosphere results in part of the outgoing terrestrial radiation being absorbed, scattered and re-emitted by the atmosphere back towards the earth. This reduces the total terrestrial flux escaping from the earth and significantly increases the earth's surface temperature,  $T_{\sigma}$ .

#### 2.4 Temperature profile

Differential absorption of energy by gas molecules within each atmospheric layer results in a vertical atmospheric temperature profile, as illustrated by the U.S. Standard Atmosphere mean vertical temperature profile shown in Fig. 3.

The heated earth emits radiation which is transmitted to the atmosphere by convective and radiative interactions. In the troposphere (convection dominated lower atmosphere) the average temperature decreases with altitude from the surface at a rate called the adiabatic lapse rate, defined by:

$$\Gamma = -dT/dz \tag{2}$$

with  $\Gamma$  in units of K km<sup>-1</sup> and z in km.

The stratosphere (radiation dominated middle atmosphere) is associated with the presence of ozone. Ozone strongly absorbs incoming ultraviolet solar radiation which heats the atmosphere and results in a gradual increase in temperature with altitude, with a peak at about 50 km. In the model, the temperature of the stratosphere is assumed to be constant at the tropopause level since the lower stratosphere is essentially isothermal (see Fig. 3) and because the pressure of the atmosphere decreases rapidly with altitude. The model temperature profile is thus taken to be:

$$T(p) = \begin{cases} T_g - \Gamma z(p) & p_T \le p \le p_g \\ T(p_T) & 0 \le p \le p_T \end{cases}$$
(3)

where  $p_{T}$  is the tropopause pressure and z(p) is the altitude at pressure p.



Figure 3. Vertical atmospheric temperature profile according to US Standard Atmosphere

#### 2.5 Pressure profile

Since the vertical motions of the average atmosphere are generally very small, the atmosphere is assumed to be in hydrostatic equilibrium. This means that each infinitesimal layer of unit cross section can be described by:

$$dp = -g\rho dz \tag{4}$$

where  $\rho$  is the atmospheric density, g is the gravitational constant and dp is the upward pressure force that is balanced by the downward force of the gas weight. Using the assumed temperature structure, the hydrostatic equilibrium equation and the ideal gas law, the altitude can be expressed as a function of pressure. The ideal gas law can be written as:

$$o = p\overline{M}/RT$$
 (5)

where  $\overline{M}$  is the mean molecular weight, R is the universal gas constant and T is the temperature. Substituting for  $\rho$  the hydrostatic equilibrium equation becomes:  $dp/p = -(g\overline{M}/RT)dz$ 

On using the model temperature profile and integrating, the altitude-pressure profile becomes:

$$z(p) = \begin{cases} T_g \left[ 1 - (p/p_g)^{\alpha} \right] / \Gamma & \text{for } p_T \le p \le p_g \\ z(p_T) - \alpha T \ln(p/p_T) / \Gamma & \text{for } 0 \le p \le p_T \end{cases}$$
(7)

with  $\alpha = \Gamma R/g\overline{M}$ .

A fixed pressure grid is chosen to divide the atmosphere into levels and this allows T(p) to be readily computed given  $T_{\alpha}$  and  $\Gamma$ .

## **3 TERRESTRIAL RADIATION**

#### 3.1 Emissivity

Throughout the surface-atmosphere system radiation is being absorbed and re-emitted. Absorbing radiation increases the kinetic energy of the particles, which at equilibrium is released as infrared radiation. The integrated wavelength flux emitted by a perfectly absorbing gas (or black body) is equal to  $\sigma T^4$ , where T is the equilibrium temperature of the body of particles and  $\sigma = 5.67 \times 10^{-8} \text{ Wm}^{-2} \text{K}^{-4}$  is the Stefan-Boltzman constant. For a non-perfectly absorbing gas or imperfect black body the emitted flux integrated over all wavelengths can be written as  $\varepsilon \sigma T^4$ , where  $\varepsilon$  is the emissivity of the body whose value ranges from 0 to 1.

The emitted flux is proportional to  $T^4$  and thus a small change in temperature produces a large change in flux. The layers of gas in the dense lower atmosphere, the ground and low clouds essentially behave as black bodies and so  $\varepsilon \simeq 1$  while the emissivity of cirrus or high clouds can be as low as 0.3. Local thermodynamic equilibrium conditions probably extend upto about the mesopause (e.g. Houghton 1979, p. 62).

#### 3.2 Optical depth

Although a body may emit a flux F at an altitude  $z(p_1)$ , only a fraction, t, of F will reach  $z(p_2)$  because of absorption in the layer between  $p_1$  and  $p_2$ . This fraction is called the transmissivity of the gas layer between  $z(p_1)$  and  $z(p_2)$ . The transmissivity can be written in terms of an 'atmospheric optical depth',  $\tau$  (see Fig. 4). If the atmospheric optical depth (or opacity) is large, then the transmissivity is low. The transmissivity between  $z(p_1)$  and  $z(p_2)$  is defined by:

$$\mathbf{t} = \exp(-\tau)$$

(8)

where  $\tau = \tau(p_1, p_2)$  is the atmospheric optical depth between  $z(p_1)$  and  $z(p_2)$ .

The wavelength integrated optical depth,  $\tau$ , is dependent on the atmospheric absorber amount, pressure and temperature. The atmospheric absorber amount, W, for each absorber, a, is the number of grams of absorber gas per unit area (cm<sup>2</sup>) between altitudes  $z(p_1)$  and  $z(p_2)$  and is given by: ŝ,

$$W_{a}(p_{1}, p_{2}) = \int_{z_{1}}^{z_{2}} \rho_{a} dz$$
 (9)

This can be expressed in terms of the mixing ratio which is given by:

$$\eta_{\rm a} = \rho_{\rm a}/\rho = p_{\rm a}M_{\rm a}/pM \tag{10}$$

where  $\rho_a$ ,  $p_a$ ,  $M_a$  are the density, pressure and molecular weight of the absorber and  $\rho$ , p, M those of air.



Figure 4. Flux transmission through an atmosphere layer between  $p_1$  and  $p_2$ 

Substituting dp/g =  $-\rho$ dz and  $\eta_a$  in equation (9), the absorber amount becomes:

$$W_{a}(p_{1},p_{2}) = - \int_{p_{1}}^{p_{2}} (\eta_{a}/g) dp, \qquad p_{1} > p_{2}$$
(11)

Although the main components of air are nitrogen, oxygen and argon, the most important absorbers are the trace constituents  $H_2O$ ,  $CH_4$ ,  $CO_2$  and  $O_3$  which strongly affect the atmospheric infrared transmission.

Water vapour decreases with altitude and this can be parameterized by:

$$\eta_{\rm H2O} = \eta_{\rm H2Og} \left(\frac{\rm p}{\rm p_g}\right)^{\beta}$$
(12)

where  $\eta_g$  refers to the surface mixing ratio at  $p_g$ .

Substituting equation (12) into (11) the water vapour absorber amount becomes:

$$W_{H2O}(p_1, p_2) = \eta_{H2Og}[p_1^{1+\beta} - p_2^{1+\beta}] / [(1+\beta)gp_g^{\beta}]$$
(13)

where  $p_1 > p_2$ .

The total water vapour amount,  $W_{H2O}^* = W_{H2O}(0,p_g)$ , is then:

$$W_{\rm H2O}^* = \eta_{\rm H2Og} \, p_g \, / \, (1 + \beta)g \tag{14}$$

Therefore the water vapour absorber amount between  $z(p_1)$  and  $z(p_2)$  in terms of  $W^*_{H2O}$  is:

$$W_{H2O}(p_1, p_2) = W_{H2O}^* \left[ (p_1/p_g)^{1+\beta} - (p_2/p_g)^{1+\beta} \right]$$
(15)

for 
$$p_1 > p_2$$
.

A value of  $\beta$  is needed in terms of meteorological data.

From equation (14):

$$\beta = \eta_{\rm H2Og} p_{\rm g} / g W_{\rm H2O}^* - 1 \tag{16}$$

Now since:

$$\mathbf{p}_{\mathbf{H}_{2\mathbf{O}}} = \mathbf{r}_{\mathbf{H}} \ \mathbf{e}_{\mathbf{s}}(\mathbf{T}) \tag{17}$$

where  $r_{\rm H}$  is the relative humidity (expressed as a fraction) and  $e_{\rm s}(T)$  is the saturation water vapour pressure (mbar) at temperature T:

$$p_{H2Og} = r_{Hg} e_s(T_g)$$
(18)

and from equation (10):

$$\eta_{\rm H2Og} = p_{\rm H2Og} \, M_{\rm H2O} / p_g \, M \tag{19}$$

Substituting the last two equations into (16) gives:

$$\beta = 0.634 r_{Hg} e_s(T_g) / W_{H2O}^* - 1$$
<sup>(20)</sup>

where  $M_{H2O}/M = 0.62$ , g = 980.665 cms<sup>-2</sup>,  $W_{H2O}^*$  is a measurable meteorological quantity representing the total vertical water vapour amount (g cm<sup>-2</sup>) and  $e_s$  in mbar is the saturation water vapour pressure at T given by:

$$\ln e_{s} = a_{1} - a_{2} / T + a_{3} \ln T$$
(21)

where  $a_1 = 58.1717$ ,  $a_2 = 6938.67$  and  $a_3 = -5.5189$  for T > 273 K.

For CO<sub>2</sub> and CH<sub>4</sub>,  $\eta_a$  can be taken to be a constant. Integrating equation (11) for CH<sub>4</sub> and CO<sub>2</sub> the absorber amounts are:

$$W_{CO2}(p_1, p_2) = W_{CO2}^* (p_1 - p_2)/p_g$$
(23)

and

$$W_{CH4}(p_1, p_2) = W_{CH4}^* (p_1 - p_2)/p_g \text{ for } p_1 > p_2$$
(24)

The ozone mixing ratio peaks within the stratosphere and the amount above level z (km) is computed from the expression of Green (1964):

$$W_{O3}(0,p) = W_{O3}^* G(z(p))$$
<sup>(25)</sup>

with  $G(z) = [1 + \exp(-b/c)] / [1 + \exp((z-b)/c)]$  (26)

where for a mid-latitude ozone distribution b = 20 km and c = 5 km are representative values (see Lacis & Hansen 1974).

The ozone absorber amount between two altitudes  $z(p_1)$  and  $z(p_2)$  is:

$$W_{O3}(p_1, p_2) = W_{O3}(0, p_1) - W_{O3}(0, p_2) \quad \text{for } p_1 > p_2$$
(27)

Substituting equation (25) into (27) gives:

$$W_{O3}(p_1, p_2) = W_{O3}^* \left[ G(z(p_1)) - G(z(p_2)) \right]$$
(28)

In Vardavas & Carver (1984b) the atmospheric infrared spectrum was divided into 22 broad intervals which have important water vapour, carbon dioxide, ozone and methane absorption bands. The flux was then computed at each atmospheric level for each spectral interval. The spectral flux was then summed to obtain the integrated flux at each atmospheric level.

In this work we are interested in computing the wavelength integrated flux at the top and at the base of the atmosphere. Since the atmospheric radiation field is primarily determined by the surface temperature  $T_g$ , it is characterised by a Planck function with a temperature in the vicinity of 200-300 K. In order to assign a wavelength averaged optical depth to the atmospheric layers we thus need to average the wavelength dependent optical depth weighted by the spectral distribution of the terrestrial radiation field and so we use the Planck mean optical depth to determine the transfer of integrated terrestrial radiation flux.

The Planck mean infrared optical depth  $\tau$  of the whole atmosphere can be expressed as:  $\tau = -\ln \overline{t}$ (29)

with 
$$\overline{t} = \int_{0}^{\infty} B_{\lambda}(T) t_{\lambda} d\lambda / \sigma T^{4} \simeq \sum_{i} \beta_{i} t_{i}$$
 (30)

where  $\beta_i(T) = \overline{B}_i(T) \Delta \lambda_i / \sigma T^4$ ,  $\overline{B}_i$  is the mean value of the Planck function within spectral interval  $\Delta \lambda_i$  and  $t_i$  is the associated mean transmission. The interval of integration extends from about 5  $\mu$ m to infinity for T  $\approx$  300 K. Following the CV model of Vardavas & Carver (1984b), the interval of integration was divided into 22 sub-intervals which have important water vapour, carbon dioxide, ozone and methane infrared absorption bands. The  $\beta_i$  were found to be weakly temperature-dependent for the range 200-300 K and it was found that a useful simplification is to set  $\beta_i(T) = \beta_i(T_*)$  with T<sub>\*</sub> = 288 K and then introduce a simple temperature correction to the optical depth  $\tau$  via:

$$c(T_g) = -0.0045T_g + 2.3$$
 with  $c(288) = 1.0$  (31)

This correction factor allows for the increase in the atmospheric mean optical depth with decreasing surface temperature. As the surface temperature decreases the peak in the Planck function, which describes the spectral flux emitted by the surface, moves to a higher wavelength and hence the relative contribution of the almost opaque water vapour bands beyond 20  $\mu$ m increases (see Fig. 5).

Using equations (15), (23), (24) and (28), the Planck mean optical depth for each absorber amount W (gm cm<sup>-2</sup>) can be readily evaluated from the following expressions based on the work of Vardavas & Carver (1984b):

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$$\begin{aligned} \tau_{\rm H2O}({\rm p_1,p_2}) &= 0.63 \ (W_{\rm H2O})^{0.52} \\ &\text{for } 0 \le W_{\rm H2O} \le 10 \end{aligned}$$

$$\begin{aligned} \tau_{\rm CO2}({\rm p_1,p_2}) &= 0.14 \ (W_{\rm CO2})^{0.22} \\ &\text{for } 0 \le W_{\rm CO2} \le 10 \end{aligned}$$

$$\begin{aligned} \tau_{\rm O3}({\rm p_1,p_2}) &= 2.51 \ (W_{\rm O3})^{0.62} \\ &\text{for } 0 \le W_{\rm O3} \le 2.10^{-3} \end{aligned}$$
and

$$\tau_{CH4}(p_1, p_2) = 2.51 (W_{CH4})^{0.75}$$
  
for  $0 \le W \le 10^{-2}$  (32)



Figure 5. The Planck function as a function of wavelength and temperature

The atmospheric optical depth between  $z(p_1)$  and  $z(p_2)$  can then be expressed as:

$$\tau(p_1, p_2) = 1.66c(T_g) \left( \tau_{H2O}(p_1, p_2) + \tau_{O3}(p_1, p_2) + \tau_{CH4}(p_1, p_2) + \tau_{CO2}(p_1, p_2) \right)$$
(33)

where 1.66 is a standard diffusivity factor introduced to allow for non-vertical radiation transfer.

Note that the expressions in equation (32) are strictly valid only for the present total atmospheric pressure.

#### 3.3 Clear sky fluxes

#### Net outgoing terrestrial flux

The net outgoing terrestrial flux consists of surface and atmospheric radiation. The surface emission is  $\varepsilon_g \sigma T_g^4$ , where  $\varepsilon_g$  is the emissivity and  $T_g$  is the temperature of the ground (see Fig. 6). However, only  $t_g \varepsilon_g \sigma T_g^4$  reaches the top of the atmosphere where  $t_g = \exp(-\tau_g)$  is the total atmospheric transmissivity,  $\tau_g = \tau(0, p_g)$  and  $p_g$  is surface pressure.

The atmosphere can be divided into pressure or optical depth layers each emitting infrared radiation according to its average gas temperature. The radiation emitted from these layers can reach the top of the atmosphere either directly or reflected from the surface. Each atmospheric layer at temperature T emits a flux equal to  $\sigma T^4$ , but only  $t_1\sigma T^4$  reaches the top of the atmosphere directly, where  $t_1 = \exp(-r(0,p))$  is the transmissivity between the layer at z(p) and the top of the atmosphere z(p = 0). Summing over all the atmospheric layers (p = 0 to  $p = p_g$ ), the flux reaching the top of the atmosphere directly is given by:

$$r_g$$
  
 $\int t_1 \sigma T^4 d\tau$ ,  $t_1 = \exp(-\tau(p,0))$  and  $r_g = \tau(0,p_g)$   
0

Each gas layer also emits  $\sigma T^4$  in the direction of the earth. This is first reduced by  $t_2 = \exp(-\tau(p,p_g))$ . Then at the surface a fraction  $\varepsilon_g$  is absorbed and a fraction  $1 - \varepsilon_g$  is reflected. Finally only  $t_g$  of the reflected flux reaches the top of the atmosphere. Therefore the flux reaching the top of the atmosphere via the earth's surface is  $t_g (1 - \varepsilon_g) t_2 \sigma T^4$ . Summing over all of the gas layers from p = 0 to  $p = p_g$ , the flux reaching the top of the atmosphere indirectly is equal to:

$$(1 - \epsilon_g) t_g \int_0^{\tau_g} t_2 \sigma T^4 d\tau, \quad t_2 = \exp(-\tau(p, p_g))$$

The net outgoing terrestrial flux under a clear sky at the top of the atmosphere is given by the sum of the above components (see Fig. 6).

$$\overline{F}_{s}^{\dagger}(\infty) = t_{g}\epsilon_{g}\sigma T_{g}^{4} + \int_{0}^{\tau g} t_{1}\sigma T^{4}d\tau + (1-\epsilon_{g})\int_{0}^{\tau g} t_{g}t_{2}\sigma T^{4}d\tau$$
(34)

The net flux can be re-written as:

$$\overline{F}_{s}^{\dagger}(\infty) = \overline{t}_{s} \varepsilon_{g} \sigma T_{g}^{4}$$
(35)

where  $\overline{t}_s$  is an effective atmospheric transmissivity for a clear sky and is given by:

$$\overline{t}_{s} = t_{g} + \frac{1}{\epsilon_{g}} \int_{0}^{\tau_{g}} \left(\frac{T}{T_{g}}\right)^{4} t_{1} d\tau$$
$$+ \frac{(1 - \epsilon_{g})}{\epsilon_{g}} t_{g} \int_{0}^{\tau_{g}} \left(\frac{T}{T_{g}}\right)^{4} t_{2} d\tau \quad (36)$$



Figure 6. The clear sky transmission of terrestrial flux emitted by the ground and each atmospheric layer

Net upward flux at the surface

Understanding the derivation of  $\overline{F}_{s}^{\dagger}(\infty)$ , it is easy to derive the equations for the net upward terrestrial flux at the surface,  $\overline{F}_{s}^{\dagger}(0)$ , a flux that is important in solar evaporation models.

At the surface, the net upward terrestrial flux for a clear sky atmosphere may be evaluated from:

$$\overline{F}_{\mathbf{s}}^{\dagger}(0) = \overline{\epsilon}_{\mathbf{s}} \sigma T_{\mathbf{g}}^{4}$$
(37)

where  $\overline{\epsilon}_{s}$  represents an effective ground emissivity for a clear sky atmosphere and is given by:

$$\overline{\epsilon}_{s} = \epsilon_{g} - \epsilon_{g} \int_{0}^{\tau_{g}} \left(\frac{T}{T_{g}}\right)^{4} t_{2} d\tau + (1 - \epsilon_{g}) \int_{0}^{\tau_{g}} \left(\frac{T}{T_{g}}\right)^{4} t_{2} d\tau$$
(38)

The first term in equation (38) is the ground emittance, the second term is the atmospheric flux absorbed by the surface, while the third is the atmospheric flux reflected by the surface (see Fig. 6).

#### 3.4 Cloud cover

Clouds cover approximately 50% of the earth and significantly affect the surface temperature by reflecting to space a fraction of the incident solar radiation. Clouds also act as thermal insulators by absorbing the upward terrestrial flux and re-radiating it back towards the ground.

The total cloud cover fraction  $A_c$  can be estimated from day length d (h) and sunshine hours from:

$$A_c = 1 - s/d \tag{39}$$

In the present model the total cloud cover, which includes the effect of overlapping cloud layers, is divided into three non-overlapping components which are contributed by low clouds, middle clouds and high clouds (see Fig. 7). The contribution  $A_{ci}$  of each cloud type to the total cloud cover is evaluated from:

$$A_{ci} = f_i e_i A_c / \sum_i f_i e_i$$
(40)

where  $f_i$  is a cloud frequency of occurrence weight and  $e_i$  is a cloud extent weight given by:

$$f_i = n_i / \sum_i n_i$$
 and  $e_i = A'_{ci} / \sum_i A'_{ci}$  (41)



Figure 7. A sky described by a clear sky and cloudy sky fraction, where the cloudy sky fraction comprises three non-overlapping layers

with  $n_i$  the frequency of occurrence of cloud type i and  $A'_{ci}$  is the observed average fractional cloud cover. The total cloud cover is then given by:

$$A_{c} = \sum_{i} A_{ci}$$
(42)

#### 3.5 Cloudy sky fluxes

#### Net outgoing terrestrial flux

The net outgoing terrestrial flux for a cloudy sky comprises the atmospheric, cloud and surface emitted flux (See Fig. 8).

Cloud emitted flux reaching the top of the atmosphere equals  $t_c \varepsilon_c \sigma T_c^4$ , where  $t_c$  is the above cloud atmospheric transmissivity,  $\varepsilon_c$  is the cloud emissivity and  $T_c$  the temperature of the cloud top, where  $t_c = \exp(-\tau(0,p_c))$  and  $p_c$  is the cloud top pressure.

From a gas layer situated above the cloud, the flux that reaches the top of the atmosphere equals  $t_1\sigma T^4$ , where  $t_1$  is the transmissivity between the layer at p and the top of the atmosphere. Summing over all the atmospheric layers above the cloud, the outgoing flux contributed by these atmospheric layers is given by:

$$\int_{0}^{\tau_{c}} t_{1}\sigma T^{4}d\tau, t_{1} = \exp(-\tau(p,0))$$

$$\int_{0}^{0} \text{ and } \tau_{c} = \tau(0,p_{c})$$



Figure 8. The four sources of outgoing terrestrial radiation for a cloudy sky are: the ground, the cloud and the atmospheric layers above and below the cloud The flux emitted from a gas layer situated below the cloud, goes through three reductions. Firstly, by transmission through the atmosphere below the cloud layer it is reduced by a factor  $t_3 = \exp(-\tau(p,p_{cb}))$ , where  $p_{cb}$  is the cloud base pressure. Secondly, the flux passes through the cloud and is reduced by a factor  $1 - \varepsilon_c$ , where  $\varepsilon_c$  is the absorptivity (equal to the emissivity) of the cloud. Thirdly, by transmission through the atmosphere above the cloud layer it is reduced by a factor  $t_c$ . Therefore the flux reaching the top of the atmosphere from a gas layer situated below the cloud is equal to  $\sigma T^4 t_3 (1 - \varepsilon_c)t_c$ . Summing over all the layers below the cloud (from  $p = p_g$  to  $p = p_{cb}$ ), the contributed flux is:

$$\tau_{cb}$$
  
(1 -  $\epsilon_c$ )t<sub>c</sub> $\int t_3\sigma T^4 d\tau$ ,  $t_3 = \exp(-\tau(p, p_{cb}))$  and  $\tau_{cb} = \tau(p_{cb}, p_g)$   
0

The surface emitted infrared flux  $\epsilon_g \sigma T_g^4$  reaching space is reduced by the atmosphere above and below the cloud and by the cloud itself. The flux reaching the top of the atmosphere from the ground is thus  $t_{cb} (1 - \epsilon_c) t_c \epsilon_g \sigma T_g^4$ , where  $t_{cb} = \exp(-\tau_{cb})$ .

The net outgoing flux under a cloudy sky,  $\overline{F}_c^{\dagger}(\infty)$ , is thus given by the sum of the above components.

$$\overline{F}_{c}^{\dagger}(\infty) = t_{c}\epsilon_{c}\sigma T_{c}^{4} + \int t_{1}\sigma T^{4}d\tau + (1 - \epsilon_{c}) t_{c} \int t_{3}\sigma T^{4}d\tau$$

$$0 \qquad 0$$

$$+ t_{cb} (1 - \epsilon_{c}) t_{c} \epsilon_{g}\sigma T_{g}^{4}$$
(43)

The net outgoing flux for a cloudy atmosphere at the top of the atmosphere can also be written as:

$$\overline{F}_{c}^{\uparrow}(\infty) = \overline{t}_{c} \ \varepsilon_{c} \sigma T_{c}^{4}$$
(44)

with

$$\overline{t}_{c} = t_{c} + \frac{1}{\epsilon_{c}} \left(\frac{Tg}{T_{c}}\right)^{4} \int_{0}^{\tau_{c}} t_{1} \left(\frac{T}{T_{g}}\right)^{4} d\tau + \frac{(1 - \epsilon_{c})}{\epsilon_{c}} t_{c} \left(\frac{Tg}{T_{c}}\right)^{4} \int_{0}^{\tau_{c}} \left(\frac{T}{T_{g}}\right)^{4} t_{3} d\tau + t_{cb} \frac{(1 - \epsilon_{c})}{\epsilon_{c}} t_{c} \epsilon_{g} \left(\frac{Tg}{T_{c}}\right)^{4}$$

$$(45)$$

The net outgoing terrestrial flux for the whole sky is the sum of clear sky and cloudy sky net outgoing fluxes weighted by the fractional cloud cover.

$$\overline{F}_{ir}^{\dagger}(\infty) = (1 - A_c) \overline{F}_s^{\dagger}(\infty) + \sum_i A_{ci} \overline{F}_c^{\dagger}(\infty)_i$$
(46)

where  $\overline{F}_{c}^{\uparrow}(\infty)_{i}$  is the net outgoing flux above cloud type i and  $\overline{F}_{s}(\infty)$  is the net outgoing flux from the clear sky.

$$\overline{F}_{ir}^{\uparrow}(\infty)$$
 can be written as:  
 $\overline{F}_{ir}^{\uparrow}(\infty) = \overline{t}_{f}\sigma T_{g}^{4}$ 
(47)

where  

$$\overline{t}_{f} = (1 - A_{c})\overline{t}_{s}\varepsilon_{g} + \sum_{i} A_{ci} \overline{t}_{ci} \varepsilon_{ci} (T_{ci}/T_{g})^{4}$$
(48)

If we let  $fr_i = A_{ci}/A_c$  then equation (46) can be written as:

$$\vec{F}_{ir}^{\uparrow}(\infty) = \Sigma fr_i \ \vec{F}_{ir}^{\uparrow}(\infty)_i$$
(49)

where

$$\overline{F}_{ir}^{\uparrow}(\infty)_{i} = (1 - A_{c})\overline{F}_{s}^{\uparrow}(\infty) + A_{c} \overline{F}_{c}^{\uparrow}(\infty)_{i}$$
(50)

In evaluating the above integrals it is important to integrate in the correct direction.

Let  $F^+$  be the total flux emitted from all layers of the atmosphere between levels  $z_1$  and  $z_2$  reaching  $p_2$  and  $F^-$  be the flux emitted from all layers of the atmosphere between levels  $z_1$  and  $z_2$  reaching  $p_1$ . Then it is important to note that  $F^+ \neq F^-$  because the density of the gas varies with altitude and so the flux reaching level  $p_1$  is dominated by radiation emitted from the adjacent layers which do not have the same absorber amounts as layers adjacent to  $p_2$ .

## Net upward flux at the surface

Below a cloud layer the net upward flux  $\overline{F}_c^{\uparrow}(0)$  has three components: the surface emittance, the absorbed flux at the surface and the atmospheric flux reflected by the surface (see Fig. 9). The surface absorbs  $\varepsilon_g$  and reflects  $(1-\varepsilon_g)$  of the incident flux which has three components. These comprise the flux from the atmosphere below the cloud layer,  $f_1$ , the flux from the cloud base,  $f_2$ , and the atmospheric flux from above the cloud layer,  $f_3$  and are given by:

$$f_1 = \int_0^{\tau_{cb}} t_2 \sigma T^4 d\tau \quad f_2 = t_{cb} \sigma T_c^4 \quad \text{and} \quad f_3 = t_{cb} (1 - \epsilon_c) \int_0^{\tau_c} t_4 \sigma T^4 d\tau \quad (51)$$

with  $t_4 = \exp(-\tau(p,p_c))$ .

The net upward flux at the surface below the cloud layer is then:

$$\overline{F}_{c}^{\uparrow}(0) = \varepsilon_{g} \sigma T_{g}^{4} + (1 - \varepsilon_{g}) (f_{1} + f_{2} + f_{3}) - \varepsilon_{g}(f_{1} + f_{2} + f_{3})$$
(52)

which can be re-written as:

$$\overline{\mathbf{F}}_{\mathbf{c}}^{\dagger}(0) = \overline{\varepsilon}_{\mathbf{c}} \ \sigma \mathbf{T}_{\mathbf{g}}^{4} \tag{53}$$

where  $\overline{\varepsilon}_{c}$  is the effective surface emissivity below the cloud given by:

$$\overline{\varepsilon}_{c} = \varepsilon_{g} + (1 - \varepsilon_{g}) (f_{1} + f_{2} + f_{3}) / \sigma T_{g}^{4}$$
$$- \varepsilon_{g} (f_{1} + f_{2} + f_{3}) / \sigma T_{g}^{4}$$
(54)

For a number of different non-overlapping cloud types i we have:



Figure 9. The emitted, absorbed and reflected components of terrestrial flux at the surface beneath a cloudy sky

$$A_{c}\overline{\epsilon}_{c} = \sum_{i} A_{ci}\overline{\epsilon}_{ci}.$$
(55)

The total net upward terrestrial flux at the ground is then:

$$\overline{F}_{ir}^{\dagger}(0) = (1 - A_c)\overline{F}_s^{\dagger}(0) + \sum_i A_{ci} \overline{F}_c^{\dagger}(0)_i$$
(56)

where  $\overline{F}_{c}^{\dagger}(0)_{i}$  is the net upward flux at the surface below cloud type i and  $\overline{F}_{s}^{\dagger}(0)$  is the net upward flux for a clear sky at the surface.

In terms of the fraction  $fr_i = A_{ci}/A_c$  equation (56) can be written as:

$$\overline{\mathbf{F}}_{i\mathbf{r}}^{\dagger}(0) = \sum_{i} \mathbf{fr}_{i} \ \overline{\mathbf{F}}_{i\mathbf{r}}^{\dagger}(0)_{i}$$
(57)

where  

$$\overline{F}_{ir}^{\dagger}(0)_{i} = (1 - A_{c}) \overline{F}_{s}^{\dagger}(0) + A_{c} \overline{F}_{c}^{\dagger}(0)_{i}$$
(58)

# **4 SOLAR RADIATION**

#### 4.1 Incoming flux

The mean daily incoming solar flux is given by:

$$F_{IN}^{\dagger} = S_{o}\overline{\mu} d/r^{2} (Jm^{-2} day^{-1})$$
 (59)

Where  $S_o$  is the solar constant (1367 Wm<sup>-2</sup>, Frohlich 1983),  $\overline{\mu}$  is the daily mean cosine of the solar zenith angle, d is the daylength (s) and  $1/r^2$  is a correction factor due to earth's elliptical orbit. The daylength is given by:

$$d = 8.64 \times 10^4 \text{ H/}\pi \quad \text{(s)} \tag{60}$$

The correction factor  $1/r^2$  can be evaluated from:

$$r = (1 - e^2)/(1 + e \cos\nu)$$
(61)

which is the ratio of the magnitude of the radius vector of the earth's orbit to the mean distance between the earth and sun, where e = 0.0167 is the orbit eccentricity,  $\nu$  is the obliquity of the sun given by  $\nu = 78^{\circ} - \lambda$  with sin  $\lambda = \sin \delta / \sin \delta_{\circ}$  where  $\lambda$  is its longitude,  $\delta$  its declination and  $\delta_{\circ} = 23.45^{\circ}$ .

The cosine of the solar zenith angle,  $\mu$ , varies throughout the day. However, the mean daily solar zenith angle  $\overline{\mu}$  is given by:

$$\overline{\mu} = A + B (\sin H)/H \tag{62}$$

where H is the hour angle given by:

$$\cos H = -\tan\phi \,\tan\delta \tag{63}$$

where  $\phi$  is the latitude of a given location on earth, and

$$A = \sin\phi \, \sin\delta \tag{64}$$

$$\mathbf{B} = \cos\phi \, \cos\delta \tag{65}$$

For seasonally and diurnally averaged conditions we can take a mean daylength of  $d = \frac{1}{2}$  day, i.e.  $H = \pi/2$ , and  $\delta = 0$  (at equinox) so that the mean solar zenith angle at latitude  $\phi$  is given by:

 $\widetilde{\mu}(\phi) = 2 \cos \phi/\pi$ 

We can then evaluate a global or latitudinal mean for the solar zenith angle by noting that the fraction of the spherical surface area contained within a latitudinal interval  $d\phi$  is  $\cos\phi d\phi$  so that the surface weighted mean is given by:

$$\overline{\mu} = \int_{0}^{\pi/2} \mu(\phi) \cos\phi d\phi$$
(67)

assuming the earth to be a perfect sphere. Thus  $\overline{\mu}$  can be evaluated from:

$$\overline{\mu} = \frac{2}{\pi} \int_{0}^{\pi/2} \cos^2 \phi d\phi$$

$$= 1/2$$
(68)

The mean global daily incoming solar flux can then be obtained from equation (59) by setting d = 1/2 day,  $\overline{\mu} = 1/2$  and  $r \simeq 1$  so that:

$$F_{1N}^{*} \simeq S_{0}/4 \quad (Wm^{-2})$$
 (69)

In order to compute the global mean surface temperature the mean global daily net incoming solar flux is needed. This can be evaluated from:

$$\mathbf{F}_{\mathbf{o}}^{\dagger}(\infty) = (1 - \mathbf{R}_{\mathbf{p}})\mathbf{F}_{\mathbf{I}\mathbf{N}}^{\dagger} = \mathbf{a} \ \mathbf{F}_{\mathbf{I}\mathbf{N}}^{\dagger}$$
(70)

where  $R_p$  is the planetary albedo and a is the planetary absorptivity.

The absorptivity, a, can be written in terms of the clear sky planetary absorptivity,  $a_s$ , and cloudy sky absorptivity  $a_{ci}$  for a cloud type i covering a fraction  $A_{ci}$  of the sky:

$$\mathbf{a} = (1 - \mathbf{A}_{c})\mathbf{a}_{g} + \sum_{i} \mathbf{A}_{ci} \mathbf{a}_{ci}$$
(71)

So that

.

$$\overline{F}_{o}^{\downarrow}(\infty) = F_{IN}^{\downarrow} \left[ (1 - A_{c})a_{s} + \sum_{i} A_{ci} a_{ci} \right]$$
(72)

This can also be written in terms of the fraction,  $fr_i = A_{ci}/A_c$ , as:

$$\overline{F}_{o}^{\downarrow}(\infty) = \sum_{i} \operatorname{fr}_{i} \overline{F}_{o}^{\downarrow}(\infty)$$

where

$$\overline{F}_{o}^{\downarrow}(\infty)_{i} = F_{IN}^{\downarrow} \left[ (1 - A_{c})a_{s} + A_{c} a_{ci} \right]$$
(73)

#### 4.2 Planetary albedo

The planetary albedo,  $R_p$ , is modelled by considering a four layer atmosphere (see Figs 10 and 11). At the top is an ozone layer which transmits only a fraction,  $t_{uv}$ , of the incoming solar flux,  $S_o/4$ , because of ultraviolet absorption between 0.2-0.35  $\mu$ m corresponding to about 4.5% of the incoming solar radiation. If all uv is absorbed, then 95.5% of incoming solar radiation will be transmitted, which corresponds to  $t_{uv}$  having a minimum value of 0.955. The ozone layer also absorbs in the visible (0.45-0.85  $\mu$ m) corresponding to 45% of

(66)



Figure 10. Clear sky solar flux transmission



Figure 11. Cloudy sky solar flux transmission

incoming solar radiation. According to the CV model of Vardavas & Carver (1984b),  $t_{uv}$  can be approximated by:

$$t_{\rm uv} = 1 - 0.023 (W_{\rm O3}^*/\bar{\mu})^{0.18}$$
(74)

and

$$t_{\rm vis} = 1 - 0.021 W_{\rm O3}^* / \overline{\mu} \tag{75}$$

where  $W_{03}^*$  is now measured in cm-STP.

For incoming downward radiation,  $\overline{\mu} = 1/2$ , as was previously calculated. However solar radiation travelling in the upward direction has been diffused by contact with either the ground or clouds and so  $\overline{\mu}$  is set to 3/5 which is the standard diffusivity approximation for converting diffuse transmission to equivalent direct transmission (e.g. see Kontratyev 1969). Both t<sub>vis</sub> or t<sub>uv</sub> agree well with fractions given in Lacis & Hansen (1974).

The second layer is a cloud layer with fractional cloud cover  $A_c$  (see Fig. 11). Each cloud type transmits a fraction  $t_{cl}$ . The simple two-stream approximation (Sagan & Pollack 1967) with no absorption is used so that  $t_{cl}$  is given by (see Lacis & Hansen 1974):

$$t_{cl} = 1 - R_{cl} \tag{76}$$

where  $R_{cl}$  is the cloud albedo given by:

$$R_{c1} = \frac{\sqrt{3(1 - g_{c1})\tau_{c1}}}{2 + \sqrt{3}(1 - g_{c1})\tau_{c1}}$$
(77)

where  $\tau_{cl}$  is the cloud optical depth to solar radiation and  $g_{cl}$  is the asymmetry factor taken to be 0.85 in the present work.

The third layer is a Rayleigh scattering layer with water vapour absorption in the near infrared. Rayleigh scattering by air is important in the uv and visible but negligible in the near-infrared. The fraction  $t_R$  is transmitted after scattering and this fraction is set to a fixed value of 0.93 (e.g. Lacis & Hansen 1974) for both upward and downward radiation. Following Lacis & Hansen, Rayleigh scattering is ignored for the atmosphere beneath a cloud layer. Approximately 40% of incoming solar radiation is in the near-infrared (> 0.85  $\mu$ m) so the minimum value of  $t_w$  is 0.60 (see Vardavas & Carver 1984b, table 6). From the complex model of Vardavas & Carver (1984),

$$t_w = 1 - 0.11 \ (W_{H2O}^*/\tilde{\mu})^{0.31}$$

using the Curtis-Godson approximation. The above expression is in close agreement with that of Yamamoto (1962).

(78)

The absorption of solar radiation by carbon dioxide takes place throughout the atmosphere. For a clear sky atmosphere the CV model gives:

$$t_{c} = 1 - 0.015 \ (W_{CO2}^{*}/\overline{\mu})^{0.263}$$
<sup>(79)</sup>

For the total atmosphere the transmissivity,  $t_c$ , is computed with an absorber amount,  $W^*_{CO2}$ . Above a cloud layer the transmission, denoted by  $t_{ca}$ , is computed for an absorber amount  $W^*_{CO2}$  ( $p_c/p_g$ ) while below a cloud layer the transmission, denoted by  $t_{cb}$ , is computed for an absorber amount  $W^*_{CO2}$  ( $1 - p_c/p_g$ ).

The last layer is, in fact, an isotropically reflecting surface with a global mean albedo,  $R_{r}$ , of about 0.1 (e.g. Henderson-Sellers & Wilson 1983).

The planetary albedo  $R_p$  may be computed from:

$$\mathbf{R}_{\mathbf{p}} = 1 - \mathbf{a} \tag{80}$$

where a is the absorption by the surface-atmosphere system:

$$\mathbf{a} = (1 - \mathbf{A}_{c})\mathbf{a}_{s} + \sum_{i} \mathbf{A}_{ci}\mathbf{a}_{ci}$$
(81)

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where  $a_s$  is the planetary absorptivity for a clear sky and  $a_{ci}$  that for a cloud type i covering a fraction  $A_{ci}$  of the sky (assuming non-overlapping cloud cover components).

The derivation of the clear sky absorptivity,  $a_s$ , is shown in Fig. 10, where F represents the incoming solar radiation at the top of the atmosphere,  $F_{IN}^{\downarrow}$ .

The three solar flux components at the top of the atmosphere are: incoming, Rayleigh scattered and flux reflected by the earth.

For a clear sky:

$$a_{s} = 1 - t_{vis}^{+} t_{uv}^{+} (1 - t_{R}) t_{vis}^{-} - R_{g} \beta^{+} \beta^{-}$$
(82)

where 
$$\beta^+ = t_R t_w^+ t_c^+ t_{vis}^+ t_u^+$$
 and  $\beta^- = t_R t_w^- t_c^- t_{vis}^-$  (83)

and the superscripts +/- refer to direct ( $\overline{\mu} = 1/2$ ) and diffuse ( $\overline{\mu} = 3/5$ ) radiation, respectively. The atmospheric transmission for upward uv flux,  $t_{uv}$ , is set to 1 since this flux component is negligible.

The deriviation of  $a_c$  is shown in Fig. 11 where below the cloud layer  $\overline{\mu} = 3/5$ . Rayleigh scattering below a cloud layer is ignored because of downward re-scattering by the cloud cover. Above a multilayer cloud cover there is no simple way for including Rayleigh scattering. The error in the total outward flux that results from omitting this term is ~ 2%.

The solar flux components at the top of the atmosphere are incoming, cloud reflected and earth reflected. The cloud reflected flux is reduced by the ozone layer and  $CO_2$  above the cloud while the earth reflected flux is reduced by ozone, cloud,  $CO_2$ , water and ground absorption.

The cloudy sky absorptivity, a<sub>c</sub>, can be thus written as:

$$a_{c} = 1 - \alpha_{c}R_{cl} - \beta_{c}R_{g}$$
(84)

where  $\alpha_c = \alpha^+ \alpha^-$  with  $\alpha^+ = t_{uv}^+ t_{vis}^+ t_{ca}^+$ ,  $\alpha^- = t_{vis}^- t_{ca}^$ and  $\beta_c = \alpha_c t_{cl}^2 (t_w^- t_{cb}^-)^2$ 

#### 4.3 Flux at surface

In evaporation models, the net daily downward solar flux  $\overline{F}_0^{\downarrow}(0)$  at the water surface is required and this can be computed, to a good approximation from:

$$\overline{F}_{o}^{\downarrow}(0) = F_{IN}^{\downarrow} \left[ (1 - A_{c})t_{g}^{*} + \sum_{i} A_{ci} t_{ci}^{*} \right]$$
(85)

where  $t_s^* = (1 - R_g) \beta^+$  and  $t_{ci}^* = (1 - R_g) \alpha^+ t_{cl} t_w^- t_{cb}^-$ 

are the effective atmospheric transmissivity for a clear sky and through a cloudy sky for cloud type i, respectively (see Figs 10 & 11). For a waterbody  $R_g \simeq 0.07$  is a reasonable estimate (e.g. Henderson-Sellers & Wilson 1983).

In terms of the fraction  $fr_i = A_{ci}/A_c$  equation (85) can be re-written as:

$$\overline{F}_{o}^{\downarrow}(0) = \sum_{i} fr_{i} \overline{F}_{o}^{\downarrow}(0)_{i}$$
(86)

where  $\overline{F}_{o}^{\downarrow}(0)_{i} = F_{IN}^{\downarrow}[(1-A_{c})t_{s}^{*} + A_{c}t_{ci}^{*}]$ 

# **5 MODEL COMPUTATIONS**

#### 5.1 Method of solution

The model can be used to determine either the mean global surface temperature of the earth and the associated atmospheric radiation fluxes at the top and bottom of the atmosphere or these fluxes alone given local meteorological conditions and a specified surface temperature.

The atmosphere is divided into fixed pressure levels which define the atmospheric layers and altitudes of the cloud layers. This allows the  $\tau$  grid structure of the atmosphere to be computed which then allows the following fluxes to be evaluated:  $\overline{F}_{o}^{\downarrow}(\infty)$ , the net downward solar flux at the top of the atmosphere;  $\overline{F}_{ir}^{\uparrow}(\infty)$ , the net upward terrestrial flux at the top of the atmosphere;  $\overline{F}_{o}^{\downarrow}(0)$ , the net downward solar flux at the surface; and  $\overline{F}_{ir}^{\downarrow}(0)$ , the net upward terrestrial flux at the surface.

Equation (1) can be solved iteratively for the surface temperature  $T_g$  by noting that  $\overline{F}_{ir}^{\uparrow}(\infty) = \overline{t}_f \sigma T_g^4$  (equation 47) where  $\overline{t}_f$  is dependent on  $T_g$ . The net incoming solar flux,  $\overline{F}_o^{\downarrow}(\infty) = a F_{IN}^{\downarrow}$ , can be taken as temperature independent so that the global mean surface temperature  $T_g$  can be obtained iteratively from:

$$T_{g} = \left[a F_{IN}^{\downarrow} / \overline{t}_{f} \sigma\right]^{\frac{1}{4}}$$
(87)

using Newton-Raphson iteration.

Note that the atmospheric optical depth due to water vapour,  $\tau_{H2O}$ , depends on the total water vapour amount  $W^*_{H2O}$  and on the vertical water vapour profile which is determined by  $\beta$  (T<sub>g</sub>) (equation 20). If  $\beta$  is arbitrarily fixed to a particular value then equation (87) can be solved rapidly since  $\overline{t_f}$  is then only weakly dependent on T<sub>g</sub>. If  $\beta$  is solved consistently with T<sub>g</sub> at each iteration then  $\overline{t_f}$  is more strongly dependent on T<sub>g</sub>.

#### 5.2 Reference atmosphere

The output of this simple model (SV) is now compared with the more complex model (CV) of Vardavas & Carver (1984b). The CV model is a one-cloud single-layer model and so for comparison the SV model is run for a single cloud model.

The global mean surface temperature and the solar and terrestrial flux variations with cloud cover, carbon dioxide and water vapour that are predicted by the SV model are found to be in very good agreement with those of the CV model and with the work of other authors.

The CV reference atmospheric model was the present global mean atmosphere described by the following parameters:

$S_o = 1367 \text{ Wm}^{-2} (F_{1N}^{\downarrow} = 342 \text{ Wm}^{-2})$	$R_g = 0.1$
$W_{CO2}^* = 0.54 \text{ g cm}^{-2} (345 \text{ ppmm})$	$W_{H2O}^* = 1.7 \text{ g cm}^{-2}$ (or precipitable cm)
$W_{O3}^* = 0.75 \times 10^{-3} \text{ g cm}^{-2} (0.35 \text{ cm}\text{-STP})$	$W_{CH4}^* = 0.9 \times 10^{-3} \text{ g cm}^{-2} (1.6 \text{ ppmm})$
$\Gamma = 6.5 \text{ K } \text{km}^{-1}$	$p_T = 200 \text{ mbar}$
$A_{c} = 0.5$	$p_c = 530 \text{ mbar}$
$\Delta p_{c} = p_{cb} - p_{c} = 20 \text{ mbar}$	$\tau_{\rm c} = 7.7$
$\varepsilon_{g} = 1$	$\varepsilon_{\rm c} = 1$

For the above parameters it can be seen from Table 1 that the SV and CV model results are in very good agreement. The net incoming solar flux at the top of the atmosphere  $\overline{F}_{o}^{4}(\infty) = 234.6 \text{ W m}^{-2}$  according to the SV reference model with  $A_{c} = 0.5$  and 234.1 W m<sup>-2</sup> according to the CV model. This is in agreement with the value of  $234 \pm 7 \text{ W m}^{-2}$  given in Stephens et al. (1981) which is based on four years of satellite observations. The asymmetry factor  $g_{cl}$  was 0.85 for both the CV and SV model and the water vapour profile generated by the CV model was found to be best matched by the SV model using  $\beta = 4$  in equation (15). For a clear sky atmosphere it can be seen from Table 1 that the two models are in close agreement.

The net downward solar flux at the surface  $\overline{F}_{0}^{\downarrow}(0)$  is 172.0 W m<sup>-2</sup> according to the SV reference model compared with 168.2 W m<sup>-2</sup> given by the CV model, 160.0 W m<sup>-2</sup> given by Peng et al. (1982) and 164.9 W m<sup>-2</sup> given by Manabe & Strickler (1964). The downward atmospheric infrared flux  $F_{1r}^{\downarrow}(0)$  at the surface is 320.7 W m<sup>-2</sup> according to the SV model compared with 317.5 W m<sup>-2</sup> for the CV model and 330.7 W m<sup>-2</sup> given by Peng et al. (1982). Note that this flux depends crucially on the value assigned to  $W_{H20}^{+2}$ .

Table 1. Comparison of SV and CV model fluxes (W m<sup>-2</sup>) and surface temperature (K)

-	$A_c = 0$		A <sub>c</sub> :	$A_c = 0.5$		$A_c = 1$	
	sv	cv	sv	cv	sv	сv	
	299.5	296.8	234.6	234.1	169.6	171.3	
<b>F</b> <sup>↓</sup> <sub>0</sub> (0)	227.9	223.1	172.0	168.2	116.1	113.5	
F <sup>↓</sup> (0)	327.3	327.4	320.7	317.5	<b>289</b> .7	285.1	
т <sub>g</sub>	296.1	296.9	287.8	288.2	275.0	275.2	

#### 5.3 Carbon dioxide effects

The variation of the global mean suface temperature  $T_g$  with atmospheric carbon dioxide (in PAL = present atmospheric level = 345 ppmm) is given in Fig. 12 with all other parameters fixed to those of the reference model. As can be seen there is a very good agreement between the CV model which predicts a temperature change  $\Delta T_g = 1.3$  K and the SV model which predicts  $\Delta T_g = 1.1$  K for a doubling of atmospheric carbon dioxide. The above changes in  $T_g$  are comparable to the lower range of values given in Climatic Change (1983) for global warming due to a doubling of the carbon dioxide and are in excellent agreement with the 1.2-1.3 K change given in Hansen et al. (1985) for a model with no climatic feedbacks. The variation of  $\Delta T_g$  with carbon dioxide will be strongly influenced by the atmospheric water vapour and cloud cover feedbacks. As was discussed in Vardavas & Carver (1985), and shown in Fig. 12, the sensitivity of  $\Delta T_g$  to a variation in carbon dioxide decreases with increasing atmospheric carbon dioxide due to the saturation of the strong carbon dioxide 15  $\mu$ m bands. The increased sensitivity for concentrations  $\leq 1$  PAL could be an important factor in triggering recurrent ice ages on earth.

The change in the downward atmospheric infrared flux at the surface  $\Delta F_{1r}^{\downarrow}(0)$  as a function of carbon dioxide amount is shown in Fig. 13. The SV and CV model results are in good agreement, both predicting  $\simeq 6 \text{ W m}^{-2}$  increase in the downward atmospheric infrared flux at the surface with doubling of CO<sub>2</sub>. However, the increase in T<sub>g</sub> results in an increase in the upward infrared flux emitted by the surface so that the net upward infrared flux at the surface decreases by  $\simeq 0.5 \text{ W m}^{-2}$ . In addition, the doubling of carbon dioxide results in a reduction in the downward solar flux at the surface by  $\simeq 0.5 \text{ W m}^{-2}$  so that the net gain of energy at the surface is close to zero.

In Fig. 14 are shown the effective ground infrared emissivities  $\overline{\epsilon}_s$  and  $\overline{\epsilon}_c$ , equations (38) and (54), corresponding to a clear sky ( $A_c = 0$ ) and a totally cloud covered ( $A_c = 1$ ) sky, respectively, as functions of the atmospheric carbon dioxide amount. Both  $\overline{\epsilon}_s$  and  $\overline{\epsilon}_c$  exhibit



Figure 12. The global mean surface temperature change  $\Delta T_g$  (K) as a function of atmospheric carbon dioxide amount in PAL (present atmospheric level = 345 ppmm) for the present simple (SV) and complex model (CV) of Vardavas & Carver (1984)



Figure 13. The change in global mean downward infrared flux (W  $m^{-2}$ ) at the surface as a function of atmospheric carbon dioxide

a weak variation with carbon dioxide. Note that the values of  $\overline{\epsilon}_s$  and  $\overline{\epsilon}_c$  in Fig. 14 are for  $\epsilon_g = 1$ , if these are denoted by  $\overline{\epsilon}_s(1)$  and  $\overline{\epsilon}_c(1)$ , then in general:

$$\overline{\varepsilon}(\varepsilon_{g}) = \varepsilon_{g} \ \overline{\varepsilon}(1) \tag{88}$$
where

 $\overline{\varepsilon}(1) = (1 - A_{c}) \overline{\varepsilon}_{s}(1) + A_{c} \overline{\varepsilon}_{c}(1)$ 

For a water surface  $\varepsilon_g$  can lie between 0.9 and 1.0 (Beard & Hollen 1969; Hobson & Williams 1971) so that setting  $\varepsilon_g = 1$  can result in a maximum error of about 10% in the net upward infrared flux.

In Fig. 14 are also shown the atmospheric transmissivities  $\overline{t}_s$  and  $\overline{t}_c$  (equations 36 & 45) for a clear sky and above a cloud layer, respectively. The SV and CV model results are once again in good agreement.

#### 5.4 Water vapor effects

In Fig. 15 the effective emissivities  $\overline{\varepsilon}_{g}$  (clear sky) and  $\overline{\varepsilon}_{c}$  (under cloud cover) computed using  $\varepsilon_{g} = 1$  are shown as functions of the atmospheric water vapour amount in g cm<sup>-2</sup> or precipitable cm, keeping all other parameters fixed to their reference model values. The emissivities show a stronger variation with water vapour than they do with carbon dioxide. The SV and CV results are in good agreement for  $W_{H2O}^* \ge 1$  cm, which is the region of interest for tropical and mid-latitudes. Atmospheric water vapour at any latitude is quite variable from one season to the next and when averaged annually it exhibits a decreasing trend with increasing latitude. In equatorial regions the mean annual value is 4.5 cm decreasing to  $\simeq 1$  cm at 60° latitude. In the wet/dry tropical northern Australia water vapour varies from  $\simeq 2$  cm during the Dry season to  $\simeq 4$  cm in the Wet.

The effective transmissivities  $\overline{t}_s$  (clear sky) and  $\overline{t}_c$  (above the cloud layer) also exhibit a stronger variation with water vapour than they do with carbon dioxide. Water vapour also affects appreciably the net incoming solar flux as can be seen from Fig. 16. The transmissivities  $t_s^*$  (clear sky) and  $t_c^*$  (total cloud cover) are shown and these describe net downward solar radiation at the surface, as defined by equation (85). The effective



Figure 14. The effective surface infrared emissivities  $\bar{\varepsilon}_{g}$  (under a clear sky) and  $\bar{\varepsilon}_{c}$  (under a cloud layer) as functions of atmospheric carbon dioxide. The effective atmospheric infrared transmissivities  $t_{g}$  (clear sky) and  $\bar{t}_{bc}$  (above cloud layer) are shown. SV\_\_\_\_, CV\_\_\_\_



Figure 15. The effective surface infrared emissivities and atmospheric infrared transmissivities as functions of atmospheric water vapour in precipitable cm. SV\_\_\_\_\_, CV\_\_\_\_\_



Figure 16. The effective planetary absorptivities  $a_g$  (clear sky) and  $a_c$  (total cloud cover) for solar radiation as functions of atmospheric water vapour (cm). The associated effective transmissivities  $t_g^*$  and  $t_c^*$  for solar radiation reaching the surface are also shown.  $SV_{---}$ ,  $CV_{---}$ 

absorptivities  $a_g$  (clear sky) and  $a_c$  (total cloud cover), which are functions of atmospheric water vapour (equations 82 and 84) and which determine the planetary albedo, are shown in Fig. 16.

The fraction f of the net downward solar flux at the surface which is lost by surface infrared emission is important to solar evaporation studies and can be written as:

$$\mathbf{f} = \overline{\mathbf{F}}_{ir}^{\dagger}(0) / \overline{\mathbf{F}}_{o}^{\downarrow}(0) \tag{89}$$

Recently, Huillet & Lauga (1985) have used a fixed value of 0.25 in their soilplant-water model to compute solar evaporation in a forested catchment. From equation (89), f will depend on various meteorological quantities, amongst them cloud cover, incoming solar flux, surface temperature, carbon dioxide, water vapour and ground emissivity. The variation of f with water vapour is shown in Fig. 17 for the three surface temperatures 0, 15 and 30°C, for a clear sky atmosphere and with all other parameters set to those of the reference model. As the water vapour increases, the downward atmospheric flux increases and hence f decreases. The variation of f with  $\Gamma$ ,  $A_c$  and carbon dioxide was found to be weak.

As  $F_{IN}^{\downarrow}$  increases f will decrease so that, given a set of meterological conditions, f should be lower at the tropics than at mid-latitudes.

#### 5.5 Latitudinal flux trends

The SV model was further tested by examining annually averaged latitudinal trends in the net downward solar flux and net upward infrared flux at the surface. The latitudinal trends were based on a model with  $R_g = 0.1$ ,  $\Gamma = 6.5$ ,  $\tau_c = 7.7$  with the  $W_{H2O}$  (cm),  $W_{O3}$  (cm-STP) and  $A_c$ trends taken from Sasamori et al. (1972) for the southern hemisphere, an adopted  $T_g$ trend based on various sources and a computed  $F_{IN}^{I}$  (W m<sup>-2</sup>), as shown in Table 2 at each latitude L (degrees). In Fig. 18 the latitudinal variation of the net downward solar flux at the surface is shown, and it can be seen that there is good agreement with the trends of other workers (London 1957; Sasamori et al. 1972; Stephens et al. 1981). In Fig. 19 the trend of the net upward infrared flux at the surface is shown for  $\varepsilon_g = 1.0$ . As can be seen the SV model results are in close agreement with the variation given by Stephens et al. (1981) for the northern hemisphere. It should be borne in mind that the SV model computations are meant to show rough latitudinal trends, especially at high latitudes.

In Fig. 20 the annually averaged latitudinal trend of f is shown for  $\varepsilon_g = 1.0$  and it can be seen that f increases from ~ 0.25 at tropical regions to ~ 0.6 at high latitudes. For artificial water reservoirs Beard & Hollen (1969) give  $\varepsilon_g = 0.922$  so that setting  $\varepsilon_g = 1$  can result in an  $\simeq 8\%$  error in f.

Table 2. Model latitudinal trends in water vapour (cm), ozone (cm-STP), cloud cover, surface temperature (K) and incoming solar flux (W  $m^{-2}$ )

L	w <sub>H2O</sub>	w <sub>O3</sub>	A <sub>c</sub>	Тg	FIN↓
0	4.2	0.26	0.48	305	415
10	3.7	0.27	0.48	300	410
20	2.9	0.28	0.48	295	394
30	2.2	0.31	0.50	290	367
40	1.6	0.33	0.59	285	331
50	1.2	0.35	0.72	280	288
60	0.8	0.35	0.83	275	<b>24</b> 1



Figure 17. The fraction f of net solar radiation at the surface which is emitted as infrared radiation as a function of water vapour and surface temperatures 0, 15 and 30 (°C)



Figure 18. The latitudinal trend of net downward solar flux (W m<sup>-2</sup>) at the surface given by the present SV model is compared with that of SCV (Stephens et al. 1981) and SLH (Sasamori et al. 1972) and L (London 1957)



Figure 19. The latitudinal trend of net upward ir flux  $(W m^{-2})$  at the surface, for  $\varepsilon_g = 1$ , given by the present SV model is compared with that given by L (London 1957) and SCV (Stephens et al. 1981)



Figure 20. The latitudinal trend of the fraction f of net emitted infrared to net absorbed solar radiation by the surface for  $\varepsilon_g=1$ 

# **6 COMPUTER PROGRAMME**

## 6.1 Programme modes and input data

The computer programme SRC (Simple Radiative-Convective) can be run in two modes: in the first mode it iteratively computes the mean global surface temperature and the solar and terrestrial fluxes at the top and bottom of the atmosphere, while in the second mode it computes the fluxes given site specific temperature conditions.

# Mode A

The first mode requires mean global atmospheric conditions as input data and the iteration flag, ITF, is set to 1. The ground temperature is solved iteratively until flux balance at the top of the atmosphere is achieved. The initial estimate for the surface temperature is  $T_0$ .

## Mode B

In the second mode, atmospheric data specific to a particular location are required, ITF is set to 0 which fixes the surface temperature at  $T_0$ . The fluxes at the top and bottom of the atmosphere are calculated for the given atmospheric data.

The input data are stored on a file called SRCIN and these are shown in the following table.

Text symbol	Computer name	Name	Units
$S_0 = 1367$	So	Solar constant	W m <sup>-2</sup>
μ	AMÜ	Mean cosine of solar zenith angle	
d	DAY	Day length	days
R <sub>σ</sub>	RG	Mean ground albedo to solar radiation	
w Ås	WO3	Atmospheric ozone amount	cm-STP
W <sup>4</sup> 20	WH2O	Atmospheric water vapour amount	g cm <sup>-2</sup>
Wču4	WCH4	Atmospheric methane amount	$g cm^{-2}$
W <sup>*</sup> CO2	WCO2	Atmospheric carbon dioxide amount	$g cm^{-2}$
ε <sub>σ</sub>	EG	Ground infrared emissivity	
<sup>e</sup> ci	ECN(N) N=1,2,3	Cloud infrared emissivity	
pc <sub>i</sub>	PCN(N)	Cloud top pressure	atmospheres
dpci	DPCN(N) N=1,2,3	Cloud thickness	mbar
$\tau_{ci}$	CLTN(N) N=1,2,3	Cloud optical depth to solar radiation	
$\alpha_{ci}$	CLAN(N)=0 N=1,2,3	Cloud absorption of solar radiation	
Pt	РТ	Pressure at tropopause	mbar
T	то	First estimate of surface temperature	ĸ
r	GLAPZ	Lapse rate in troposphere	K km <sup>-1</sup>
<sup>p</sup> g	PG	Pressure at ground	mbar

#### INPUT PARAMETERS

Text symbol	Computer name	Name	Units
A <sub>c</sub>	cc	Fractional cloud cover	
f <sub>ri</sub>	FRN(N) N=1,2,3	Fraction of total cloud cover due to cloud type i	$A_{ci}/A_{c}$
<sup>r</sup> Hg	RHO	Relative humidity at ground	
-	ITF	Iteration flag	

# INPUT PARAMETERS (contd)

## 6.2 Flow diagram

A flow diagram of the computer programme is given in Fig. 21. The programme first reads the input data file SRCIN then sets up the atmospheric pressure grid which defines the atmospheric layers and the location of the cloud layers, as shown in Fig. 22. The corresponding altitude grid, for a given atmospheric temperature structure, is then computed by a call to ZET. The programme then computes the terrestrial and solar fluxes for each cloud type by calls to subroutines TERRE and SOLAR, respectively. To compute the terrestrial fluxes, the programme needs to call three subroutines. Subroutine TAUS sets up the atmospheric optical depth grid, corresponding to the pressure grid. Subroutines HTAU and HINT then compute the integrals associated with the terrestrial flux transfer. Once the fluxes at the top and base of the atmosphere are known, the iteration flag ITF determines whether a global mean surface temperature will be computed iteratively via the Newton-Raphson temperature correction method. If ITF = 0 the programme run is for local atmospheric conditions and the programme exits via a call to ALTI which generates a table of the vertical atmospheric temperature and optical depth structure. If ITF = 1 the programme is automatically run once again with entry point the subroutine ZET which computes the altitude grid for a given atmospheric temperature structure. When the surface temperature, and hence temperature structure, converges (usually within about 5 iterations) the programme exits via a call to ALTI. The output is written on the file SRCOUT.

A detailed description of each programme subroutine now follows.

# 6.3 Main programme

The main programme first reads the input data and calculates  $F_{IN}^{\downarrow}$ . The parameters for the Newton-Raphson calculation of  $T_g$  are defined and shown in the following table. The main programme calls PGRID to set up the pressure grid. The temperature and altitude profiles are set up for the present value of  $T_g$  by calling subroutine ZET. The solar and terrestrial flux under each cloud type, both at the top and at the base of the atmosphere, are computed by calls to subroutines TERRE and SOLAR. The computed fluxes are:  $\overline{F}_{ir}^{\uparrow}(0)_i$ ,  $\overline{F}_{ir}^{\uparrow}(\infty)_i$ ,  $\overline{F}_0^{\downarrow}(0)_i$ ,  $\overline{F}_0^{\downarrow}(\infty)_i$  and are denoted in the programme by FIRG, FLOUT, FSOLG and FLIN, respectively. These fluxes are then summed over the 3 cloud types to give  $\overline{F}_{ir}^{\uparrow}(0)$ ,  $\overline{F}_{ir}^{\uparrow}(\infty)$ ,  $\overline{F}_0^{\downarrow}(0)$  and  $\overline{F}_0^{\downarrow}(\infty)$ , denoted in the programme by sum5, sum2, sum4 and sum1, respectively.

At this point the programme can take three possible paths:

i) If ITF = 1 and if ER (see following table) is less than  $\varepsilon$ , then the net incoming solar flux effectively equals the outgoing terrestrial flux at the top of the atmosphere (i.e. radiative equilibrium). In this case the iterations stop and  $T_g$  is the global mean surface temperature.



Figure 21. Flow diagram of the computer programme



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Figure 22. The pressure grid set up by subroutine PGRID

ii) If ITF = 1 and ER is greater than  $\varepsilon$ , then T<sub>g</sub> needs to be corrected by  $\Delta T$  until ER is less than  $\varepsilon$ . The temperature correction,  $\Delta T$ , is calculated using the Newton-Raphson method:

$$\Delta T = \frac{G(T)}{G'(T)} , \text{ where } G(T) = \begin{vmatrix} -\frac{1}{F_0} & -\frac{1}{F_1} \\ -\frac{1}{F_0} & -\frac{1}{F_1} \\ -\frac{1}{F_1} -\frac{1}{F_1} \\$$

However G(T) is a complicated function of temperature T. So G'(T) is approximated by:

$$G'(T_k) = \frac{G(T_k) - G(T_{k-1})}{T_k - T_{k-1}}$$

where  $T_k$  is the ground temperature of the present iteration and  $T_{k-1}$  is that of the previous iteration. This gives at the kth iteration a temperature correction:

$$\Delta T_{k} = \frac{G(T_{k}) (T_{k} - T_{k-1})}{G(T_{k}) - G(T_{k-1})}$$

The temperature used in the next iteration is given by  $T_{k+1} = T_k + \Delta T_k$ . For k = 1 there is no value for  $T_{k-1}$  and  $\Delta T_1$  is arbitrarily set to 10.

The programme keeps iterating until ER <  $\varepsilon$ . In the programme G(T<sub>k</sub>), G(T<sub>k-1</sub>), T<sub>k</sub>, and T<sub>k-1</sub> are denoted by GA, GB, TG and TGB, respectively.

iii) If ITF = 0 then the surface temperature is fixed at  $T_0$ . The solar and terrestrial fluxes at the top and bottom of the atmosphere are computed for this surface temperature.

Text symbol	Computer name	Name	Units
F <sub>IN</sub> ↓	FLUX	Mean daily incoming solar flux at top of atmosphere	W m <sup>-2</sup>
ε	EP	Maximum relative error allowed in flux balance equation	
	NMAX	Maximum number of iterations allowed	
	к	Iteration counter	
	ICB	Grid index for base of cloud	
	IC	Grid index for top of cloud	
$\overline{F}_{0}^{\downarrow}(\infty)_{i}$	FLIN	Net incoming solar flux at top of the atmosphere above cloud cover i	W m <sup>-2</sup>
$\overline{\mathbf{F}}_{ir}^{\uparrow}(\infty)_{i}$	FLOUT	Net outgoing terrestrial flux at top of the atmosphere above cloud cover i	W m <sup>-2</sup>
$\overline{F}_{0}^{\downarrow}(0)_{i}$	FSOLG	Net downward solar flux at the ground below cloud cover i	$W m^{-2}$
$\overline{F}_{ir}^{\dagger}(0)_{i}$	FIRG	Net upward terrestrial flux at the ground below cloud cover i	$W m^{-2}$
$\overline{F}_{0}^{\downarrow}(\infty)$	SUM1	Net incoming solar flux at top of the atmosphere	$W m^{-2}$
$\overline{F}_{ir}^{\uparrow}(\infty)$	SUM2	Net outgoing terrestrial flux at top of the atmosphere	W m <sup>-2</sup>
$\overline{F}_{0}^{\downarrow}(0)$	SUM4	Net downward solar flux at ground	W m <sup>-2</sup>
$\overline{\mathbf{F}}_{ir}^{\dagger}(0)$	SUM5	Net upward terrestrial flux at ground	$W m^{-2}$
G(T <sub>k</sub> )	GA	$\left  \overline{F}_{o}^{\downarrow}(\infty) - \overline{F}_{ir}^{\uparrow}(\infty) \right $ at iteration k.	$W m^{-2}$

#### MAIN PROGRAMME PARAMETERS

#### MAIN PROGRAMME PARAMETERS (contd)

Text symbol	Computer name	Name	Units
G(T <sub>k-1</sub> )	GB	$ \overline{F}_{o}^{\downarrow}(\infty) - \overline{F}_{ir}^{\uparrow}(\infty) $ at iteration k-1.	W m <sup>-2</sup>
т <sub>к</sub>	ТG	Ground temperature at iteration k	
т к-1	TGB	Ground temperature at iteration k-1	
	ER	$\mathbf{ER} = \left  \left( \overline{\mathbf{F}}_{ir}^{\uparrow}(\infty) - \overline{\mathbf{F}}_{o}^{\downarrow}(\infty) \right) / \overline{\mathbf{F}}_{o}^{\downarrow}(\infty) \right $	
e <sub>s</sub>	ES	Saturated water vapour pressure at ground	
β	BETA	Water vapour profile parameter (equation 20)	

# 6.4 Subroutines

#### Subroutine PGRID

Subroutine PGRID sets up a pressure grid and locates the cloud and tropopause levels. Between the ground  $(p = p_g)$  and the top of the atmosphere (p = 0) there are 100 layers of 10 mbar each. Each interval I, corresponds to a pressure p(I). At the ground I = 1, and at the top of the atmosphere I=IMAX=101. (See Fig. 22).

ICN(N) locates the top of cloud type N, where the cloud top pressure is P(ICN(N)) and N = 1,2,3 for three cloud types. Similarly, the interval at the base of cloud type N is called ICBN (N), with a pressure of p(ICBN(N)). The subroutine locates the tropopause with a pressure  $p(IT) = p_T$ .

Subroutine	PGRID	parameters
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Computer name	Name	Units
DP	Pressure grid interval set equal to 10 mbar	mbar
ICN(N)	Level corresponding to the top of cloud type N	
PC=P(ICN(N))	Pressure at top of cloud type N	atmospheres
DPC	Cloud width	mbar
PCB=PC+DPC(N	) Pressure at bottom of cloud type N Note: PCB > PC	atmospheres
ICBN(N)	Level corresponding to base of cloud type l	N

#### Subroutine ZET

ZET sets up the temperature versus pressure profile (equation 3) and the altitude versus pressure profile (equation 7).  $T_g$ ,  $\Gamma$ ,  $P_T$  and  $P_g$  are all given input parameters. The tropopause altitude corresponding to  $p = p_T$  is calculated from equation (7).

Subroutine	ZET	parameters
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Text symbol	Computer name	Name	Units
Г	GLAPZ	Lapse rate = - $dT/dz$	$K \ km^{-1}$
g	GO=9.81	Gravitational acceleration	m s <sup>-2</sup>
M	AMW=28.97	Mean molecular weight of dry air	$g  \mathrm{mol}^{-1}$
R	R=8.314	Universal gas constant	$\mathbf{J} \text{ mol}^{-1} \mathbf{K}^{-1}$
α	ALFA	$\alpha = \Gamma R/g\overline{M}$	-
$Z(p_t)$	ZT	Tropopause altitude	km

	~ .		
Text symbol	Computer name	Name	Units
p <sub>t</sub>	РТ	Tropopause pressure	mbar
	PI	P(I) mbar	
	PR	P(I)/PG -	
	IT	Tropopause level	-
	ТТ	Tropopause temperature	к

# Subroutine TERRE

Subroutine ZET parameters (contd)

TERRE computes the net terrestrial flux at the top and base of the atmosphere by first evaluating  $\overline{t}_c$ ,  $\overline{t}_s$ ,  $\overline{\varepsilon}_s$  and  $\overline{\varepsilon}_c$  using subroutines HINT and TAUS. TAUS evaluates  $\tau_c$ ,  $\tau_g$  and  $\tau_{cb}$  for  $t_c$ ,  $t_g$  and  $t_{cb}$ , while HINT evaluates the integrals.  $\overline{F}_s^{\dagger}(\infty)$ ,  $\overline{F}_c^{\dagger}(\infty)$ ,  $\overline{F}_s^{\dagger}(0)$  and  $\overline{F}_c^{\dagger}(0)$  are then found using equations 35, 44, 37 and 53. Finally  $\overline{F}_{ir}^{\dagger}(\infty)_i$  and  $\overline{F}_{ir}^{\dagger}(0)_i$  are evaluated using equations 48 and 58.

Subroutine TERRE parameters

Text symbol	Computer name	Name	Units
σ	$SIG = 5.6696 \times 10^{-8}$	Stefan-Boltzman constant	W m <sup>-2</sup> K <sup>-4</sup>
	TC	Temperature of cloud top	к
	тсв	Temperature of cloud base	к
$\tau_{g}$	TAUG	$\tau(\mathbf{p}_{\mathbf{g}},0)$	
τ <sub>c</sub>	TAUC	$\tau(\mathbf{p_c}, 0)$	
<sup>7</sup> cb	TAUCB	$\tau(\mathbf{p_{g}},\mathbf{p_{cb}})$	
tg	TRG	Total atmospheric infrared transmissivity	
t <sub>c</sub>	TRC	Atmospheric transmissivity above cloud layer	
t <sub>cb</sub>	TRCB	Atmospheric transmissivity below cloud layer	
	FSOUT	Value of the first integral in $\overline{t_s}$	
	FSO	Value of the second integral in $\overline{t}_s$ or the integrals in $\overline{\epsilon}_s$	
	FCOUT & FCUB	Value of the first and second integrals, respectively, in $\overline{t_c}$	
	FCDB & FCDA	Value of the integrals in the $f_1$ and $f_2$ terms of $\overline{\epsilon_c}$	
t <sub>s</sub>	TBS	Effective atmospheric transmissivity for a clear sky (equation 36)	
$\overline{\mathrm{F}}_{\mathrm{s}}^{\dagger}(\infty)$	FCL	$\overline{\mathrm{t}}_{\mathrm{s}}\varepsilon_{\mathrm{g}}\sigma\mathrm{T}_{\mathrm{g}}^{\mathrm{4}}$	$W m^{-2}$
<del>t</del> <sub>ci</sub>	TBC	Effective atmospheric transmissivity for a cloudy sky (equation 45)	
$\overline{F}_{c}^{\dagger}(\infty)_{i}$	FAC	$\overline{\mathrm{t}}_{\mathrm{ci}}\varepsilon_{\mathrm{ci}}\sigma\mathrm{T}_{\mathrm{c}}^{4}$	W m <sup>-2</sup>
$\overline{F}_{ir}^{\uparrow}(\infty)_i$	FLOUT	Net outgoing terrestrial flux at the top of atmosphere for cloud type i (equation 47)	W m <sup>-2</sup>
$\overline{\epsilon}_{c}$	EBC	Effective ground emissivity below a cloud layer	$W m^{-2}$
$\overline{\epsilon}_{s}$	EBS	Effective ground emissivity for a clear sky	
$\overline{F}^{\uparrow}_{1}(0)$	FBS	Net upward terrestrial flux at surface below a clear sky	W m <sup>-2</sup>

# Subroutine TERRE parameters (contd)

Text symbol	Computer name	Name	Units
$\overline{\mathbf{F}}_{\mathrm{ir}}^{\dagger}(0)_{\mathrm{i}}$	FIRG	Net upward terrestrial flux at surface below cloud i (equation 58)	<b>W</b> m <sup>-2</sup>
	FBSD	Clear sky downward ir flux absorbed by ground	
	FBSU	Upward ir flux from the ground for a clear sky atmosphere. It is the sum of flux emitted by the ground plus terrestrial flux reflected by the ground	
	FBCD	Downward ir flux absorbed by ground below a cloudy sky atmosphere	
	FBCU	Upward ir flux from the ground for a cloudy sky. It is the sum of flux emitted from the ground plus terrestrial flux reflected by the ground	

## Subroutine TAUS

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TAUS determines the atmospheric optical depth, TAU, between pressure levels p(I1) and p(I2).

First the atmospheric water, carbon dioxide, methane and ozone amounts are found between the two specified pressure levels p(I1) and p(I2). For the evaluation of  $W_{O3}^*$  subroutine OZ is called. Finally the optical depths for the four types of molecules and the TAU are calculated using equations 32 and 33.

Text symbol	Computer name	Name	Units
c(T <sub>0</sub> )	FT	Optical depth temperature correction factor	
$G(z_1)$	OZF1	Green's function at $z_1$ (equation 26)	
$G(z_2)$	OZF2	Green's function at $z_2$	
WH2O	WH2O	Total atmospheric water amount	$g cm^{-2}$
w <sup>*</sup> <sub>CO2</sub>	W C02	Total atmospheric carbon dioxide amount	$g cm^{-2}$
w <sub>03</sub>	WO3	Total atmospheric ozone amount	cm-STP
W <sup>*</sup> <sub>CH4</sub>	WCH4	Total atmospheric methane amount	g cm <sup>-2</sup>
$W_{H2O}(p_1, p_2)$	WH	Water amount between $p_1$ and $p_2$ (equations 15)	g cm <sup>-2</sup>
$W_{CO2}(p_1,p_2)$	wc	$CO_2$ amount between $p_1$ and $p_2$ (equation 23)	g cm <sup>-2</sup>
$W_{CH4}(p_1,p_2)$	WCH	Methane amount between $p_1$ and $p_2$ (equation 24)	$g cm^{-2}$
$W_{O3}(p_1,p_2)$	wo	Ozone amount between $p_1$ and $p_2$ (equation 28)	cm-STP
$\tau_{\rm H2O}(p_1,p_2)$	TAH	Optical depth of water between $\mathbf{p}_1$ and $\mathbf{p}_2$	
$\tau_{\rm CO2}({\rm p_{1},p_{2}})$	TAC	$\mathrm{CO}_2$ optical depth between $\mathtt{p}_1$ and $\mathtt{p}_2$	
$\tau_{O3}(p_1, p_2)$	TAO	${ m O}_3$ optical depth between ${ m p}_1$ and ${ m p}_2$	
$\tau_{CH4}(p_1, p_2)$	TACH	$\operatorname{CH}_4$ optical depth between $\operatorname{p}_1$ and $\operatorname{p}_2$	
$\tau(p_1,p_2)$	TAU	The total optical depth between $p_1$ and $p_2$	

Subroutine TAUS parameters

# Subroutine OZ

OZ computes Green's function G(Z) from equation (26). OZF is the computer name for G(z).

#### Subroutine HTAU

HTAU gives the integrand, H( $\tau$ ), in the equations for  $\overline{t}_s$ ,  $\overline{t}_c$ ,  $\overline{\varepsilon}_s$  and  $\overline{\varepsilon}_c$  which takes the form:

$$H(\tau) = \left(\frac{T}{T_g}\right)^4 \exp(-\tau)$$

 $\tau$  is specified by the subroutine HINT as  $\tau(P(IP),P)$ .

Subroutine	HTAU	parameters
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Text symbol	Computer name	Name	Units
$H(\tau)$	НР	Integrand in flux equations	
τ	TAU	$\tau(P,P(IP))$	
$T(\tau)$	Т	Temperature at P	

# Subroutine HINT

HINT evaluates the integrals, FIP, in the  $\overline{t}_s$ ,  $\overline{t}_c$ ,  $\overline{\epsilon}_s$  and  $\overline{\epsilon}_c$  equations. The integrals are evaluated using the trapezoidal rule. FIP represents the integral of  $H(\tau)$  between layers defined by P(I1) and P(IM) evaluated for  $\tau(P(IP), P(I))$  where I ranges from IM to I1.

For example the second term of the  $\overline{t}_{s}$  equation contains the integral:

$$= \int_{0}^{\tau} g\left(\frac{T}{T_g}\right)^4 \exp\left(-\tau(0,p)\right) d\tau$$
$$= \int_{0}^{\tau} H(\tau) d\tau, \quad \tau = \tau(0,p)$$

Therefore,

$$\tau(P(IP),P) = \tau(0,p)$$
  
$$\tau(P(IP),P(I1)) = \tau(0,p_g) = \tau_g$$
  
$$\tau(P(IP),P(IM)) = 0$$

which gives IP = IMAX, II = I and IM = IMAX. This integral is evaluated in the programme by calling HINT (IMAX, 1, IMAX, FSOUT).

## Subroutine SOLAR

SOLAR computes the net solar flux at the top and base of the atmosphere by first calculating atmospheric transmissivities. The cloud transmissivity is calculated with equation 76 and 77. The transmissivities of each absorber are solved using equations 74, 75, 78 and 79, for a given value of AMU =  $\overline{\mu}^+$  for direct radiation. For diffuse radiation  $\overline{\mu}^-$  is set equal to 3/5 whose inverse is 1.66. 'P' and 'M' as the last letter in the computer names refers to  $\overline{\mu}$  equal to  $\overline{\mu}^+$  and  $\overline{\mu}^-$ , respectively. The net solar flux at top and base of the atmosphere are determined from equation 73, 86. The computer names for the terms used are:

Text Symbol	Computer Name	Name
R <sub>CL</sub>	RC	Cloud reflectivity
$^{t}CL$	CTR	Cloud transmissivity
$\alpha_{CL} = 0$	CLA	Cloud absorptivity (set to zero in this model)
$\overline{\mu}^+$		Direct radiation mean direction
$\overline{\mu}^{-}$		Diffuse radiation mean direction
tw <sup>+</sup>	TWP	Water vapour transmissivity for direct radiation
<sup>t</sup> w <sup>-</sup>	TWM	Water vapour transmissivity for diffused radiation
t <sup>+</sup> vis	TVP	Visible transmissivity for direct radiation
t <sup>-</sup> vis	TVM	Visible transmissivity for diffused radiation
<sup>t</sup> uv	TUV	Ultraviolet transmissivity for direct downward radiation
<sup>t</sup> R	$\mathbf{TR} = 0.93$	Rayleigh scattering transmissivity (for both upward and downward radiation)
t <sub>ca</sub> +	TCPP	t <sub>CO2</sub> above cloudy sky for direct radiation
t <sub>ca</sub> -	TCPM	$t_{CO2}$ above cloudy sky for diffused radiation
tc+	TCP	$t_{CO2}$ for direct radiation for whole sky
t <sub>c</sub> -	TCM	$t_{CO2}$ for diffused radiation for whole sky
t <sub>cb</sub> "	TCPD	$t_{ m CO2}$ for diffused radiation below a cloudy sky
α <sup>+</sup>	АР	<sup>t</sup> uv <sup>t</sup> vis <sup>t</sup> ca
α_	AM	$t_{vis}t_{ca}$
β <sup>+</sup>	BP	${}^{t}\mathbf{R}{}^{t}\mathbf{w}{}^{t}\mathbf{c}{}^{t}\mathbf{v}{}^{t}_{is}{}^{t}{}_{uv}$
$\beta^-$	BM	${}^{t}\mathbf{R}^{t}\mathbf{\tilde{v}}^{t}\mathbf{\tilde{c}}^{t}\mathbf{\tilde{v}}_{is}$
α <sub>c</sub>	AA	$\alpha^+\alpha^-$
$\beta_{g}$	BB	$\beta^+\beta^-$
twtca	BC	
β <sub>c</sub>	BD	$\alpha_{\rm c} t_{\rm cl}^{\ 2} (t_{\rm w}^{-} t_{\rm cb}^{-})^2$
a <sub>s</sub>	AS	planetary absorptivity for a clear sky
<sup>a</sup> ci	AC	planetary absorptivity for cloud type i
$a_{ci}F_{IN}$	FSC	Net incoming solar flux at top of the atmosphere for a cloudy sky for cloud type i
a <sub>s</sub> F <sub>IN</sub> ↓	FSNC	Net incoming solar flux at top of the atmosphere for a clear sky
$\overline{F}_{o}^{\downarrow}(\infty)_{i}$	FLIN	Net incoming solar flux at top of atmosphere for cloud type i
<b>F</b> <sup>↓</sup> <sub>o</sub> (0) <sub>i</sub>	FSOLG	Net downward solar flux at ground below cloud type i

Subroutine SOLAR parameters

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# Subroutine ALTI

ALTI writes the altitude, pressure, temperature,  $\tau(p, p_g)$ ,  $\tau(p, p_{cb})$ ,  $\tau(p, p_c)$  and  $\tau(p,0)$  for  $0 \le z \le 30$  km and z = 100 km (for  $z \rightarrow \infty$ ). These are described as z, p, T, TAP, TAUCB, TAC and TAUD in the programme.

# 6.5 Computer code

PROGRAM SRC

```
COMMON /MT1/FLUX,CC,RG,EG
   COMMON /MT2/TG, TT, GLAPZ, PT, PG
   COMMON /MT3/IC, ICB, PC, DPC, CLT, CLA, EC, TC
   DIMENSION FRN(3), PCN(3), DPCN(3), CLTN(3), CLAN(3), ECN(3)
   DIMENSION ICN(3), ICBN(3)
   COMMON /SPEC/WCO2,WH20,WO3,WCH4,BETA
   DATA SIG/5.6696E-8/
   OPEN(2, FILE='SRCOUT', STATUS='UNKNOWN')
   OPEN(3, FILE='SRCIN', STATUS='UNKNOWN')
   READ(3,23) SO, AMU, DAY, RG, EG, PG
  C, GLAPZ, PT, CC
  C, WCO2, WH2O, WO3, WCH4, RHO
  C, (PCN(N), N=1, 3)
 C, (DPCN(N), N=1, 3)
  C, (CLTN(N), N=1, 3)
  C, (CLAN(N), N=1, 3)
  C, (ECN(N), N=1, 3)
  C, (FRN(N), N=1, 3)
  C,TO,ITF
23 FORMAT(33(8X, F9.4, /), 8X, I2)
   FLUX=AMU*DAY*S0
   WRITE(2,19)
   WRITE(2,20)
19 FORMAT(1H1,50X,'SIMPLE CLIMATE MODEL',///)
20 FORMAT(4X,'SO',6X,'AMU',3X,'DAY',4X,'RG',3X,'FLUX',
  C5X, 'CC', 3X, 'WH2O', 3X, 'WCO2', 4X, 'WO3',
  C3X, 'WCH4', 5X, 'LAP', 3X, 'PT', 7X, 'EG', 3X, 'RH0', /)
   WRITE(2,10) SO,AMU, DAY, RG, FLUX, CC, WH2O, WCO2, WO3,
  CWCH4, GLAPZ, PT, EG, RHO
10 FORMAT(2X, F7.2, 3(2X, F4.2), 2X, F7.2, 2X, F4.2,
  C3(2X,F5.2),2X,F7.5,2X,F4.1,2X,F4.0,2(3X,F4.2))
   WRITE(2,21)
   WRITE(2,22) (PCN(N), N=1,3), (DPCN(N), N=1,3), (CLTN(N), N=1,3),
  C(CLAN(N), N=1, 3), (ECN(N), N=1, 3), (FRN(N), N=1, 3)
21 FORMAT(//,10X,'CLOUD PROPERTIES'//,9X,'LOW',5X,'MIDDLE',3X,'HIGH')
22 FORMAT(//, 3X, 'PC', 3(4X, F4.2)/2X, 'DPC', 3(4X, F4.1)/2X, 'CLT',
  C3(4X,F4.1)/2X,'CLA',3(4X,F4.1)/3X,'EC',3(4X,F4.2)/3X,'FR',
  C3(4X, F4.2)///)
   EP=0.0001
   NMAX=20
   TG=T0
   CALL PGRID(PCN, DPCN, ICN, ICBN)
   K=0
 1 CONTINUE
   K=K+1
   WRITE(2,*)
```

```
WRITE(2,*) ' ITERATION = ',K
 ALES-58.1717-6938.67/TG-5.5189*ALOG(TG)
 ES-EXP(ALES)
 BETA=0.634*RH0*ES/WH20-1
 WRITE(2,*) ' TG = ',TG,' BETA = ',BETA
 CALL ZET
 SUM1=0
 SUM2=0
 SUM3=0
 SUM4=0
 SUM5=0
 DO 4 N=1,3
 WRITE(2, *)
 IF(N.EQ.1) WRITE(2,*) ' LOW CLOUD'
 IF(N.EQ.2) WRITE(2,*) ' MIDDLE CLOUD'
 IF(N.EQ.3) WRITE(2,*) ' HIGH CLOUD'
 ICB=ICBN(N)
 IC=ICN(N)
 PC=PCN(N)*PG
 DPC=DPCN(N)
 CLT=CLTN(N)
 CLA=CLAN(N)
 EC = ECN(N)
 CALL SOLAR (AMU, FLIN, FSOLG, TSTARS, TSTARC, AS, AC)
 CALL TERRE(ITF, FLOUT, FIRG, TBS, TBC, EBS, EBC)
  SUM1=SUM1+FRN(N)*FLIN
  SUM2=SUM2+FRN(N)*FLOUT
 SUM4=SUM4+FRN(N)*FSOLG
  SUM5=SUM5+FRN(N)*FIRG
4 CONTINUE
  GA=SUM2-SUM1
  DT = -GA * (TG - TGB) / (GA - GB)
  TGB-TG
  IF(K.EQ.1) TG=TGB+10
  IF(K.GT.1) TG=TGB+DT
  GB=GA
  IF(ITF.EQ.0) TG=T0
  ER=ABS((SUM2-SUM1)/SUM1)
  WRITE(2,*) K,SUM1,SUM2,DT,TG,ER
  WRITE(2,*) ' NET IR OUT = ',SUM2,' NET SOL IN = ',SUM1
  WRITE(2,*) ' NET SOLG DOWN = ',SUM4,' NET IRG UP = ',SUM5
  EGST4=EG*SIG*TG**4
  WRITE(2,*) ' GROUND EMISSION = ', EGST4
  IF(K.LT.NMAX.AND.ER.GT.EP.AND.ITF.EQ.1) GO TO 1
  IF(ER.GT.EP.AND.ITF.EQ.1) WRITE(2,*) ' NO CONVERGENCE '
5 CONTINUE
  IF(ITF.EQ.1) WRITE(2,*) ' ITF=1: SURFACE TEMPERATURE COMPUTED'
  IF(ITF.EQ.0) WRITE(2,*) ' ITF=0: SURFACE TEMPERATURE SPECIFIED'
  CALL ALTI
  STOP
  END
```

SUBROUTINE TERRE(ITF,FLOUT,FIRG,TBS,TBC,EBS,EBC)

COMMON /SPEC/WCO2,WH20,WO3,WCH4,BETA COMMON /MT1/FLUX,CC,RG,EG

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COMMON /MT2/TG, TT, GLAPZ, PT, PG COMMON /MT3/IC, ICB, PC, DPC, CLT, CLA, EC, TC COMMON /PTC/P(101), T(101), Z(101), IMAX, IT DATA SIG/5.6696E-8/ TC=T(IC)CALL TAUS(1, IMAX, TAUG) CALL TAUS(IC, IMAX, TAUC) CALL TAUS(1, ICB, TAUCB) TRG=EXP(-TAUG) TRC=EXP(-TAUC) TRCB-EXP(-TAUCB) T4 = TG \* \* 4TC4-TC\*\*4 WRITE(2,\*) TG,TC,TT,TAUG,TAUC CALL HINT(IMAX, 1, IMAX, FSOUT) CALL HINT(IMAX, IC, IMAX, FCOUT) CALL HINT(1,1,IMAX,FSO) CALL HINT(ICB, 1, ICB, FCUB) CALL HINT(1,1,ICB,FCDB) CALL HINT(IC, IC, IMAX, FCDA) EGST4=EG\*SIG\*T4 TBS=TRG+FSOUT/EG-(1-EG)\*TRG\*FSO/EG FCL=TBS\*EGST4 TBC=TRC+(T4/TC4)\*FCOUT/EC+(1-EC)\*(FCUB+EG\*TRCB)\*TRC\*T4/(EC\*TC4) FAC=TBC\*EC\*SIG\*TC4 FLOUT=(1-CC)\*FCL+CC\*FAC WRITE(2,\*) TBS,TBC WRITE(2,\*) ' IROUT CLEAR = ', FCL, ' IROUT CLOUD = ', FAC WRITE(2,\*) ' IROUT TOTAL = ', FLOUT FBSD=EGST4\*(-FSO) FBSU=EGST4+(1-EG)\*FBSD/EG FBS=FBSU-FBSD FBCD=EGST4\*(-FCDB+EC\*(TC4/T4)\*TRCB-(1-EC)\*FCDA\*trcb) FBCU=EGST4+(1-EG)\*FBCD/EG FBC=FBCU-FBCD EBC=FBC/(SIG\*T4) EBS=FBS/(SIG\*T4) WRITE(2,\*) EBS,EBC FIRG=(1-CC)\*FBS+CC\*FBC WRITE(2,\*) FBS,FB,CC,FIRG RETURN END

# SUBROUTINE SOLAR(AMU,FLIN,FSOLG,TSTARS,TSTARC,AS,AC)

COMMON /SPEC/WCO2,WH2O,WO3,WCH4,BETA COMMON /MT1/FLUX,CC,RG,EG COMMON /MT2/TG,TT,GLAPZ,PT,PG COMMON /MT3/IC,ICB,PC,DPC,CLT,CLA,EC,TC DATA SIG/5.6696E-8/,TR/0.93/ DATA GS/0.85/,R3/1.732/ PCG-PC/PG R3G=R3\*(1-GS)\*CLT RC=R3G/(2+R3G) CTR=1-RC-CLA TWP=1-.11\*(WH2O/AMU)\*\*0.31

TWM=1-.11\*(1.66\*WH20)\*\*0.31 TVP=1-.021\*W03/AMU TVM=1-.021\*1.66\*W03 TUV=1-0.023\*(WO3/AMU)\*\*0.18TCPP=1-0.015\*(WCO2\*PCG/AMU)\*\*0.263 TCPM=1-0.015\*(WC02\*PCG\*1.66)\*\*0.263 TCP=1-0.015\*(WCO2/AMU)\*\*0.263 TCM=1-0.015\*(1.66\*WCO2)\*\*0.263 TCPD=1-0.015\*(WCO2\*(1-PCG)\*1.66)\*\*0.263 AP-TUV\*TVP\*TCPP AM=TVM\*TCPM BP=TR\*TWP\*TCP\*TVP\*TUV BM-TR\*TWM\*TCM\*TVM AA=AP\*AM BB=BP\*BM BC=TWM\*TCPD BD-AA\*CTR\*CTR\*BC\*BC AC=1-AA\*RC-BD\*RG FSC=FLUX\*AC AS=1-(1-TR)\*TUV\*TVP\*TVM-BB\*RGFSNC-FLUX\*AS FLIN=CC\*FSC+(1-CC)\*FSNC TSTARS = BP\*(1-RG)TSTARC = AP \* CTR \* BC \* (1 - RG)FSOLG=(1-CC)\*FLUX\*TSTARS+CC\*FLUX\*TSTARC RETURN END

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# SUBROUTINE HINT(IP,I1,IM,FIP)

DIMENSION TAUP(101) COMMON /PTC/P(101), T(101), Z(101), IMAX, IT DO 2 I=1, IMAX CALL TAUS(IP, I, TAU) TAUP(I)=TAU 2 CONTINUE I2=I1+1IM1=IM-1 DP=TAUP(I2) - TAUP(I1) TP=T(I1)TAUPP=TAUP(I1) CALL HTAU(TP, TAUPP, HP) SUM1 = -DP\*HP/2DO 1 I=I2,IM1 DP=TAUP(I+1)-TAUP(I-1) TP = T(I)TAUPP=TAUP(I) CALL HTAU(TP, TAUPP, HP) SUM1=SUM1-DP\*HP/2 **1** CONTINUE DP=TAUP(IM)-TAUP(IM1) TP=T(IM)TAUPP=TAUP(IM) CALL HTAU(TP, TAUPP, HP)

```
FIP=SUM1-DP*HP/2
WRITE(2,*) ' FIP =',FIP
RETURN
END
```

SUBROUTINE PGRID(PCN,DPCN,ICN,ICBN)

```
DIMENSION PCN(3), DPCN(3), ICN(3), ICBN(3)
  COMMON /PTC/P(101), T(101), Z(101), IMAX, IT
 COMMON /MT2/TG, TT, GLAPZ, PT, PG
  IMAX=101
  DP=PG/(IMAX-1)
  P(1)=PG
 DO 1 I-1, IMAX
  P(I) = PG - DP*(I - 1)
  DO 2 N=1,3
  PC=PCN(N)*PG
  PCB=PC+DPCN(N)
  IF(P(I).EQ.PC) ICN(N)=I
  IF(P(I).EQ.PCB) ICBN(N)=I
2 CONTINUE
  IF(P(I).EQ.PT) IT=I
1 CONTINUE
  RETURN
  END
  SUBROUTINE TAUS(11,12,TAU)
  COMMON /SPEC/WCO2,WH2O,WO3,WCH4,BETA
  COMMON /PTC/P(101), T(101), Z(101), IMAX, IT
  P1=P(I1)
  P2=P(I2)
  PG=P(1)
  Z1=Z(I1)
  Z_{2=Z(I_{2})}
  CALL OZ(Z1,OZF1)
  CALL OZ(Z2,OZF2)
  FT = -0.0045 * T(1) + 2.3
  PR=ABS((P1-P2)/PG)
  PR5=(P1/PG)**(1+BETA)-(P2/PG)**(1+BETA)
  PR5=ABS(PR5)
  WH=WH2O*PR5
  WC=WCO2*PR
  WO-0,00214*W03*ABS(0ZF1-0ZF2)
  WCH=WCH4*PR
  TAH=0.63*WH**0.52
  TAC=0.14*WC**0.22
  TAO=2.51*WO**0.62
  TACH=2.51*WCH**0.75
  TAU=1.66*FT*(TAH+TAC+TAO+TACH)
  RETURN
  END
```

# SUBROUTINE ZET

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```
COMMON /MT2/TG,TT,GLAPZ,PT,PG
 COMMON /PTC/P(101), T(101), Z(101), IMAX, IT
 DATA G0/9.80665/,AMW/28.97/,R/8.314/
 ALFA=GLAPZ*R/(GO*AMW)
  ZT=TG*(1-(PT/PG)**ALFA)/GLAPZ
  IMAX1=IMAX-1
  DO 1 I=1, IMAX1
  PI=P(1)
  PR=PI/PG
  IF(PI.GE.PT) Z(1)=TG*(1-PR**ALFA)/GLAPZ
  T(I) = TG - GLAPZ \times Z(I)
  IF(PI.LT.PT) T(I)=TG-GLAPZ*ZT
  IF(PI.LT.PT) Z(I)=ZT-ALOG(PI/PT)*ALFA*T(I)/GLAPZ
  IF(Z(I).EQ.ZT) IT=I
1 CONTINUE
  TT=T(IT)
  Z(IMAX)=100.
  T(IMAX) = TG - GLAPZ * ZT
  RETURN
  END
```

# SUBROUTINE OZ(Z,OZF)

DATA B/20./,C/5./ AN=1+EXP(-B/C) DEN=1+EXP((Z-B)/C) OZF=AN/DEN RETURN END

# SUBROUTINE HTAU(T,TAU,HP)

COMMON /MT2/TG,TT,GLAPZ,PT,PG HP=EXP(-TAU)\*((T/TG)\*\*4) RETURN END

### SUBROUTINE ALTI

COMMON /PTC/P(101),T(101),Z(101),IMAX,IT COMMON /MT3/IC,ICB,PC,DPC,CLT,CLA,EC,TC WRITE(2,19) 19 FORMAT(////,20X,'VERTICAL STRUCTURE',//) WRITE(2,20) 20 FORMAT(4X,'Z',7X,'P',5X,' T ',6X,'TAP',5X,'TACB', C4X,' TAC',3X,' TAUD') D0 1 I=1,IMAX CALL TAUS(1,I,TAP) CALL TAUS(1,ICB,TACB) CALL TAUS(1,IC,TAC) CALL TAUS(I,IMAX,TAUD) CALL HTAU(T(I),TAP,HP)

```
CALL HTAU(T(I),TAUD,HM)

FT--0.0045*T(1)+2.3

30 FORMAT(5(2X,F6.4))

WRITE(2,21) Z(I),P(I),T(I),TAP,TACB,TAC,TAUD

21 FORMAT(2X,F4.1,2(2X,F7.2),4(2X,F6.4))

1 CONTINUE

RETURN

END
```

## 6.6 Samples of input/output

Sample input data files are given as: SRCINA (global mean) and SRCINB (specific location). For SRCINA, ITF = 1 and  $T_o = 270$  K. As can be seen in the output file SRCOUTA the program converges to the mean global temperature ( $T_g = 288.46$  K) in 5 iterations and the net inward solar flux at the top of the atmosphere equals 228 Wm<sup>-2</sup>. For SCRINB, ITF = 0 and  $T_o$  is fixed at 288.5 K. The fluxes  $\overline{F}_{ir}^{\uparrow}(\infty)$ ,  $\overline{F}_{0}^{\downarrow}(\infty)$ ,  $\overline{F}_{ir}^{\uparrow}(0)$  and  $\overline{F}_{0}^{\downarrow}(0)$  are computed for both cases. The corresponding output files are SRCOUTA and SRCOUTB.

### Input for a 3 cloud model, with ITF = 1 Global average conditions Filename = SRCINA

So	=	1367.0		solar constant (Wm <sup>-2</sup> )
Mu	=	0.5		$\overline{\mu}$ = daily mean cosine of solar zenith angle
Day	=	0.5		daylength in days
Rg	=	0.13		ground albedo
Eg	=	1.00		ground infrared emissivity
Pg	=	1000.		ground atmospheric pressure
Gam	=	6.5		lapse rate = $\Gamma = -dT/dz$ (K km <sup>-1</sup> )
Pt	=	200.		pressure of tropopause (mbar)
Ac	=	0.54		fractional cloud cover
CO2	=	0.54		$W^*_{CO2}$ = total atmospheric CO <sub>2</sub> amount (gm cm <sup>-2</sup> )
H2O	=	2.17		$W_{H_{2O}}^{*}$ = total atmospheric H <sub>2</sub> O amount (gm cm <sup>-2</sup> )
O3	=	.31		$W_{O3}^{*}$ = total atmospheric O <sub>3</sub> amount (cm-STP)
CH4	=	.0009		$W_{CH4}^{*}$ = total atmospheric $CH_{4}$ amount (gm cm <sup>-2</sup> )
rh0	=	0.77		relative humidity at surface
Pcl	=	0.76	)	
Pcm	=	0.56	Ł	pressure height at cloud top (atmospheres)
Pch	æ	0.20	J	
Dpcl	=	100.	Ì	
Dpcm	=	50.	ł	cloud thickness in mbar
Dpch	=	70.	}	
T¢l	=	20.0	ĺ	
Tcm	=	15.0	ł	cloud optical depth to solar radiation
Tch	=	2.0	J	
Tal	=	0.	Ì	
Tam	=	0.	ł	cloud absorption of solar radiation is ignored
Tah	=	0.	J	
Ecl	=	1.0	Ì	
Ecm	=	1.0	ł	cloud emissivity in infrared
Ech	=	0.3	J	-

Frl Frm Frh To	= = =	0.34 0.28 0.38 270.0	}	cloud type fra surface temper	ctional cover	uess (K)
ITF	Ŧ	1		iteration flag	ITF = 1: ITF = 0:	solves for surface temperature specified temperature

Output for a 3 cloud model, with ITF = 1 Global average conditions Filename = SRCOUTA

# SIMPLE CLIMATE MODEL

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S0	AMU	DAY	RG	FLUX	сс	WH2O	WCO2	WO3	WCH4	LAP	РТ	EG	RHO
1367.	.50	.50	.13	341.75	.54	2.17	.54	.31	.00090	6.5	200.	1.00	.77
CLOUI	LOUD PROPERTIES												
	1	LOW	М	IIDDLE		HIGH							

PC	.76	. 56	. 20
DPC	100.	50.0	70.0
CLT	20.0	15.0	2.0
CLA	0.0	0.0	0.0
EC	1.00	1.00	. 30
FR	. 34	. 28	. 38

ITERATION = 1 TG = 270. BETA = 8.76284E-2

LOW CLOUD

TG	TC	$\mathbf{TT}$	TAUG	TAUC
270.	256.27	198.8	1.9897	1.7281
FIP = .32794				
FIP = .29111				
FIP =7403				
FIP = .50152				
FIP =53044	ŀ			
FIP =56617	7			
TBS	TBC			
.46468	.53633			
IROUT CLEAR -	40,01	IROUT CLOUD	= 131.15	
IROUT TOTAL -	135.22			
EBS	EBC			
.2597	.09970			

CC FBS FBC FIRG 78.249 30.041 .54 52.216 MIDDLE CLOUD TGTCTTTAUG TAUC 270. 241.8 198.8 1,9897 1,4797 FIP = .32794FIP = .25489FIP = -.7403FIP = .54682FIP = -.6718FIP = -.41922TBC TBS .46468 .62395 IROUT CLEAR = 140.01 IROUT CLOUD = 120.94IROUT TOTAL = 129.71EBS EBC .2597 .14764 CC FIRG FBS FBC 44.484 .54 60.016 78.249 HIGH CLOUD TAUC TC TTTAUG ΤG 198.8 1.9897 .88926 198.8 270. FIP = .32794FIP = .17337FIP = -.7403 FIP = .41462 FIP = -.7266FIP = -.17341TBS TBC 4.323 .46468 IROUT CLEAR = 140.01 IROUT CLOUD = 114.84IROUT TOTAL = 126.42EBS EBC .23534 .2597 CC FIRG FBS FBC 70.909 .54 74.286 78.249 K SUM1 SUM2 DT ΤG ER .42808 1 227.89 -270. 280. 130.33 NET IR OUT = 130.33 NET SOL IN = 227.89NET SOLG DOWN = 162.21 NET IRG UP = 62.787GROUND EMISSION = 348.49ITERATION = 2TG = 280. BETA = 1.228 LOW CLOUD TAUG TAUC ΤG  $\mathbf{TC}$ TT280. 265.76 206.16 1,9072 1.4471 FIP - .3975 FIP = .31942FIP = -.77368FIP = .58FIP = -.61506FIP = -.55318

TBS TBC .546 .62883 IROUT CLEAR = 190.27 IROUT CLOUD = 177.84IROUT TOTAL - 183.56 EBS EBC .22632 .08386 CC FBS FBC FIRG 29,226 78,869 .54 52,062 MIDDLE CLOUD  $\mathbf{TG}$ TC ΤT TAUG TAUC 280. 250.76 206.16 1.9072 1.0772 FIP = .3975FIP = .2432FIP = -.77368FIP = .59523FIP = -,73623FIP = -.37252TBS TBC .546 .71861 IROUT CLEAR = 190.27 IROUT CLOUD = 161.09IROUT TOTAL = 174.52 EBS EBC .12104 .22632 FBS FBC CC FIRG 78,869 42.179 .54 59.057 HIGH CLOUD TC TGTTTAUG TAUC 280. 206.16 206.16 1.9072 .44938 FIP = .3975FIP = .1064FIP = -.77368FIP ⇒ .45855 FIP = -.76884FIP = -.10647TBS TBC . 546 5. IROUT CLEAR = 190.27 IROUT CLOUD = 153.62IROUT TOTAL = 170.48EBS EBC .22632 .20443 CC FBS FBC FIRG . 54 74.749 78.869 71.24 K SUM1 SUM2 DT TGER 2 227.89 176.06 291.34 11.335 .22743 NET IR OUT = 176.06 NET SOL IN = 227.89 NET SOLG DOWN = 162.21 NET IRG UP = 62.641GROUND EMISSION = 408, 44ITERATION = 3TG = 291.34 BETA = 3.6938LOW CLOUD ΤG TC TTTAUG TAUC 1.8136 291.34 276.52 214.5 1.0421 FIP = .49661

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FIP = .32579FIP = -.79221FIP = .65187FIP = -.69444FIP = -.48122TBS TBC .65968 .75414 IROUT CLEAR = 269.44 IROUT CLOUD = 249.97IROUT TOTAL = 258.92EBS EBC .20779 .06696 FBS FBC CC FIRG 84.87 27.347 .54 53.807 MIDDLE CLOUD ΤG TC TTTAUG TAUC 291.34 260.91 214.5 1.8136 .60971 FIP = .49661FIP = .18684FIP - .79221 FIP = .63132FIP = -.77721 FIP = -.26137TBS TBC .83396 .65968 IROUT CLEAR = 269.44 IROUT CLOUD = 219.12IROUT TOTAL = 242.26EBS EBC .20779 9.85295E-2 FBS FBC CC FIRG 84.87 40.243 . 54 60.771 HIGH CLOUD TGTCTTTAUG TAUC 291.34 214.5 214.5 1.8136 .21765 FIP = .49661FIP = 5.74964E-2FIP = -.79221FIP = .53169FIP = -.78953FIP = -5.74989E-2TBS TBC .65968 5.9506 IROUT CLEAR = 269.44 IROUT CLOUD = 214.28IROUT TOTAL = 239.65EBS EBC .20779 .18839 FIRG FBS FBC CC 76.943 80.589 84.87 . 54 SUM2 K SUM1 DT TGER 3 227.89 246,93 -3.0464 288.29 .083592 NET IR OUT = 246.93 NET SOL IN = 227.89NET SOLG DOWN = 162.21 NET IRG UP = 65.934GROUND EMISSION = 391.62ITERATION = 4TG = 288.29 BETA = 2.8674

LOW CLOUD TTΤG TCTAUG TAUC 273.63 212.26 1.8388 288.29 1.1571 FIP - .47081 FIP = .33091FIP = -.78895FIP = .63497FIP = -.67549FIP = -.50755TBS TBC .62983 .72215 IROUT CLEAR = 246.65 IROUT CLOUD = 229.51IROUT TOTAL = 237.4EBS EBC .21105 .07130 FBS FBC CC FIRG 82.652 27,924 .54 53.099 MIDDLE CLOUD TG ΤС TTTAUG TAUC 212.26 288.29 258.18 1.8388 .72599 FIP = .47081FIP = .20782FIP = -.78895FIP = .623FIP = -.76948FIP = -.2958TBS TBC .8069 .62983 IROUT CLEAR = 246.65 IROUT CLOUD = 203.28IROUT TOTAL = 223.23EBS EBC .1037 .21105 FBS FBC CC FIRG 82.652 40.612 .54 59.95 HIGH CLOUD ΤG TC ΤT TAUG TAUC 288.29 212.26 212.26 1.8388 .25155 FIP = .47081FIP = 6.53759E-2FIP = -.78895FIP = .51021FIP = -.78611FIP = -6.53858E-2TBC TBS 5.7089 .62983 IROUT CLEAR = 246.65 IROUT CLOUD = 197.11IROUT TOTAL = 219.9EBS EBC .21105 ,19133 CC FBS FBC FIRG 82.652 74.93 .54 78.482 K SUM1 SUM2 DT ΤG ER 288.46 4 227.89 226.78 .16713 .004852 NET IR OUT = 226.78 NET SOL IN = 227.89

NET SOLG DOWN = 162.21 NET IRG UP = 64.663GROUND EMISSION = 392.53 ITERATION = 5TG = 288.46 BETA = 2.9092LOW CLOUD ΤT TC TAUG ΤG TAUC 212.38 288.46 273.78 1.8374 1.1509 FIP = .47227FIP = .33079FIP = -.78916FIP = .63595FIP - .67658 FIP = -.50624TBS TBC .6315 .72397 IROUT CLEAR = 247.88 IROUT CLOUD = 230.62IROUT TOTAL = 238.56EBS EBC .21084 .07106 FBS FBC CC FIRG 27.893 53.133 82.762 .54 MIDDLE CLOUD ΤG TC TTTAUG TAUC 212.38 288.46 1.8374 .71933 258.33 FIP = .47227FIP = .20677FIP = -.78916FIP = .62348FIP = -.76997FIP = -.29396TBS TBC .8085 .6315 IROUT CLEAR = 247.88 IROUT CLOUD = 204.15**IROUT TOTAL = 224.27** EBS EBC .1034 .21084 CC FBC FIRG FBS 82.762 40.588 .54 59.988 HIGH CLOUD TC TTTAUG TAUC ΤG .24924 288.46 212.38 212.38 1.8374 FIP = .47227FIP = 6.48471E-2FIP = -.78916FIP = .51138FIP = -.78633FIP = -6.48564E-2TBS TBC 5.7227 .6315 IROUT CLEAR = 247.88 IROUT CLOUD = 198.04IROUT TOTAL = 220.97EBS EBC .21084 .19115

FBS	FBC	CC	FIRG				
82.762	75.032	. 54	78.588				
K SUM1	SUM2	DT	TG	ER			
5 227.89	227.87	.002	288.46	.000047			
NET IR OUT	= 227.87 NE	T SOL IN =	= 227.89				
NET SOLG DOWN = 162.21 NET IRG UP = 64.725							
GROUND EMISSION = 392.54							
ITF=1: SUR	FACE TEMPERA	TURE COMPU	JTED				

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# VERTICAL STRUCTURE

Z	Р	Т	TAP	TACB	TAC	TAUD
0.0	1000.00	288.46	.0000	1.7803	1.7916	1.8374
.1	990.00	287.90	.3631	1.7477	1.7591	1.8051
. 2	980.00	287.35	.4979	1.7154	1.7270	1.7731
. 3	970.00	286.79	. 5989	1.6835	1.6952	1.7414
. 3	960.00	286.22	.6823	1.6519	1.6637	1,7100
.4	950.00	285.65	.7541	1.6206	1.6325	1,6790
. 5	940.00	285.08	.8177	1.5896	1,6016	1,6483
.6	930.00	284.50	.8749	1,5590	1.5711	1.6179
.7	920.00	283.92	. 9270	1.5286	1,5409	1,5879
. 8	910.00	283.33	.9748	1.4986	1.5110	1,5581
. 9	900.00	282.73	1.0190	1.4689	1.4814	1,5287
1.0	890.00	282.13	1.0601	1.4395	1.4522	1.4997
1.1	880.00	281.53	1.0984	1.4104	1.4233	1.4709
1.2	870.00	280.91	1.1344	1.3816	1.3946	1.4425
1.3	860.00	280.30	1.1681	1.3532	1.3663	1.4143
1.4	850.00	279.67	1.1999	1.3250	1.3384	1.3865
1.4	840.00	279.05	1.2298	1.2972	1.3107	1.3591
1.5	830.00	278.41	1.2582	1.2697	1.2834	1.3319
1.6	820.00	277.77	1.2850	1.2424	1,2563	1.3051
1.7	810.00	277.12	1.3104	1,2155	1,2296	1.2786
1.8	800.00	276.47	1.3345	1,1889	1.2032	1.2524
1.9	790.00	275.81	1.3574	1.1626	1. <b>1</b> 771	1.2265
2.0	780.00	275.14	1.3792	1.1367	1.1513	1.2010
2.2	770.00	274.47	1.3999	1.1110	1.1259	1.1758
2.3	760.00	273.78	1.4196	1.0856	1.1007	1.1509
2.4	750.00	273.09	1.4383	1.0605	1.0759	1.1263
2.5	740.00	272.40	1.4562	1.0358	1.0514	1.1020
2.6	730.00	271.69	1.4732	1.0113	1.0272	1.0781
2.7	720.00	270.98	1.4895	.9871	1,0033	1.0544
2.8	710.00	270.26	1.5049	.9633	.9797	1.0311
2.9	700.00	269.53	1.5197	.9397	. 9564	1.0081
3.0	690.00	268.80	1,5338	.9164	, 9334	.9854
3.1	680.00	268.05	1.5472	, 8934	.9107	.9631
3.3	670.00	267.30	1,5600	.8708	. 8884	.9410
3.4	660.00	266.53	1.5722	.8484	,8663	.9193
3.5	650.00	265.76	1.5839	.8262	.8446	.8979
3.6	640.00	264.98	1.5950	.8044	.8231	.8768
3.7	630.00	264.19	1.6056	.7829	.8020	.8560
3.9	620.00	263.38	1.6157	.7616	.7811	.8355
4.0	610.00	262.57	1.6253	.7407	.7605	.8154
4.1	600.00	261.75	1.6345	.7199	.7403	.7956
4.2	590.00	260.91	1.6433	. 6995	.7203	.7760
4.4	580.00	260.06	1.6516	.6794	. 7007	.7568
4.5	570.00	259.20	1.6595	.6595	.6813	.7379

4,6	560.00	258,33	1.6671	.6398	.6622	.7193
4.8	550.00	257.45	1.6743	.6205	.6434	.7011
4.9	540.00	256.55	1.6812	.6013	.6249	.6831
5.0	530.00	255.64	1.6878	.5825	.6067	.6654
5.2	520.00	254.72	1.6940	. 5638	. 5888	.6481
5.3	510.00	253.78	1.6999	.5454	.5711	.6310
5.5	500,00	252.82	1.7055	. 5273	.5537	.6143
5.6	490.00	251.85	1.7109	. 5093	. 5366	. 5979
5.8	480,00	250.87	1.7160	.4916	.5198	.5818
5.9	470.00	249.87	1.7208	.4741	.5032	.5659
6.1	460.00	248.85	1.7254	.4567	.4869	. 5504
6.3	450.00	247.81	1.7298	.4395	.4708	. 5352
6.4	440.00	246.75	1.7340	.4225	.4550	. 5203
6.6	430.00	245.67	1.7379	.4057	.4394	. 5057
6.8	420.00	244.58	1.7417	.3889	.4241	.4914
6.9	410.00	243.46	1.7452	. 3722	. 4090	.4774
7.1	400.00	242.32	1.7486	.3556	. 3942	.4637
7.3	390.00	241.15	1.7518	.3391	.3795	.4503
7.5	380.00	239.96	1.7549	. 3225	.3650	.4372
7.6	370.00	238.75	1.7578	. 3058	.3508	.4244
7.8	360.00	237.51	1.7605	.2890	.3367	.4119
8.0	350.00	236.24	1.7632	. 2720	. 3228	.3996
8.2	340.00	234.94	1.7657	.2546	. 3090	.3877
8.4	330.00	233.61	1.7680	.2366	. 2953	.3760
8.6	320.00	232.25	1.7703	.2178	.2817	. 3646
8.9	310.00	230.85	1.7725	.1978	.2681	. 3535
9.1	300.00	229.41	1.7745	.1758	.2545	. 3427
9.3	290.00	227.94	1.7765	.1504	.2409	. 3322
9.5	280.00	226.42	1.7784	.1174	.2272	.3219
9.8	270.00	224.86	1.7803	.0000	.2132	.3120
10.0	260.00	223.25	1.7820	.1152	.1988	. 3022
10.3	250.00	221.59	1.7837	.1442	.1837	. 2928
10.6	240.00	219.88	1./854	.1643	.16/8	,2836
10.8	230.00	218.11	1./8/0	.1800	.1502	. 2746
	220.00	216.27	1.7885	. 1929	.1299	.2659
11.4	210.00	214.36	1.7901	.2038	.1032	.2575
11./	200.00	212.38	1.7916	.2132	.0000	.2492
12.0	190.00	212.38	1.7930	.2214	.1014	.2412
12.4	180.00	212.38	1.7945	.2286	.1248	. 2335
12./	1/0.00	212.38	1.7960	.2351	.1408	.2259
13.1	150.00	212,38	1.7975	.2410	.1532	.2185
13.5	150.00	212.38	1.7990	. 2464	.1634	.2113
13.9	140.00	212.38	1.8006	. 2513	.1/20	.2042
14.4	130.00	212.38	1.8022	.2559	.1795	.1972
14.9	120.00	212.38	1.8038	.2602	. 1862	.1904
15.4	110.00	212.38	1.8055	.2642	. 1922	.1836
16.0	100.00	212.38	1.8074	. 2682	. 1977	.1/68
10./	90.00	212,38	1,8093	. 2720	.2029	.1/00
1/.4	80.00	212.38	1.8114	.2758	.2079	. 1630
18,Z	/0.00	212.38	1.813/	.2/9/	. 212/	.1558
19.Z	60.00	212,38	1.8162	.2836	.21/5	.1481
20.3	50,00	212.38	1.8189	. 28/8	.2224	.1399
∠⊥./ )) F	40,00	212.38	1.0220	. 2921	. 22/4	.1306
23,2	30.00	212.38 010 00	1 0000	. 2900	. 2320	.1190
20,0	20.00	212.30 010 00	1.0293 1.0295	. JUTQ	1062.	. 1028
JU, J 100	10.00	212,38	1.0333 1.037/	. 30/1	. 2439	.0000
TUO'	0,00	ZIZ,30	1,83/4	, 3120	. 2492	0.0000

Input for a 3 cloud model, with ITF = 0 Local conditions Filename = SRCINB

$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$			
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	So	-	1367.0
$\begin{array}{rcl} \text{Day} = & 0.5\\ \text{Rg} = & 0.13\\ \text{Eg} = & 1.00\\ \text{Pg} = & 1000.\\ \text{Gam} = & 6.5\\ \text{Pt} = & 200.\\ \text{Ac} = & 0.54\\ \text{CO2} = & 0.54\\ \text{H2O} = & 2.17\\ \text{O3} = & .31\\ \text{CH4} = & .0009\\ \text{rh0} = & 0.76\\ \text{Pcm} = & 0.56\\ \text{Pch} = & 0.20\\ \text{Dpc} = & 100.\\ \text{Dpcm} = & 50.\\ \text{Dpch} = & 70.\\ \text{Tcl} = & 20.0\\ \text{Tcl} = & 20.0\\ \text{Tch} = & 2.0\\ \text{Tah} = & 0.\\ \text{Tah} = & 0.\\ \text{Tah} = & 0.\\ \text{Ecl} = & 1.0\\ \text{Ech} = & 0.34\\ \text{Frm} = & 0.28\\ \text{Frh} = & 0.38\\ \text{To} = & 288.5\\ \text{OTE} = & 0.38\\ \text{To} = & 288.5\\ \text{TTE} = & 0.38\\ \text{To} = & 288\\ \text{To} = $	Mu	-	0.5
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	Day	-	0.5
Eg = $1.00$ Pg = $1000$ .Gam = $6.5$ Pt = $200$ .Ac = $0.54$ CO2 = $0.54$ H2O = $2.17$ O3 = $.31$ CH4 = $.0009$ rh0 = $0.77$ Pcl = $0.76$ Pcm = $0.56$ Pch = $0.20$ Dpc = $100$ .Dpcm = $50$ .Dpch = $70$ .Tcl = $20.0$ Tch = $2.0$ Tat = $0$ .Tat = $0$ .Tat = $0$ .Ecl = $1.0$ Ech = $0.33$ Fr1 = $0.34$ Frm = $0.28$ Frh = $0.38$ To = $288.5$ TTE = $0.38$	Rg		0.13
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	Eg	-	1.00
$\begin{array}{rcl} {\rm Gam} = & 6.5 \\ {\rm Pt} = & 200 \\ {\rm Ac} = & 0.54 \\ {\rm CO2} = & 0.54 \\ {\rm H2O} = & 2.17 \\ {\rm O3} = & .31 \\ {\rm CH4} = & .0009 \\ {\rm rh0} = & 0.76 \\ {\rm Pcn} = & 0.76 \\ {\rm Pcn} = & 0.76 \\ {\rm Pcn} = & 0.56 \\ {\rm Pch} = & 0.20 \\ {\rm Dpc} = & 100 \\ {\rm Dpc} = & 100 \\ {\rm Dpc} = & 100 \\ {\rm Dpch} = & 70 \\ {\rm Tcl} = & 20.0 \\ {\rm Tcl} = & 20.0 \\ {\rm Tch} = & 2.0 \\ {\rm Tal} = & 0 \\ {\rm Tal} = & 0 \\ {\rm Tah} = & 0 \\ {\rm Cam} = & 1.0 \\ {\rm Ecl} = & 1.0 \\ {\rm Ecl} = & 1.0 \\ {\rm Ech} = & 0.38 \\ {\rm Frh} = & 0.38 \\ {\rm To} = & 288.5 \\ {\rm ITE} = & 0 \\ {\rm Opc} = & 0 \\ {\rm Opc} = & 0 \\ {\rm ITE} = & 0 \\ {\rm Opc} = & 0 \\ {\rm Opc} = & 0.38 \\ {\rm ITE} = & 0 \\ {\rm Opc} = & 0 \\ {\rm Opc} = & 0 \\ {\rm Opc} = & 0.38 \\ {\rm ITE} = & 0 \\ {\rm Opc} = & 0 \\ {$	Pg	-	1000.
$\begin{array}{rcrcrc} {\rm Pt} &=& 200.\\ {\rm Ac} &=& 0.54\\ {\rm CO2} &=& 0.54\\ {\rm H2O} &=& 2.17\\ {\rm O3} &=& .31\\ {\rm CH4} &=& .0009\\ {\rm rh0} &=& 0.76\\ {\rm Pch} &=& 0.76\\ {\rm Pch} &=& 0.76\\ {\rm Pch} &=& 0.20\\ {\rm Dpc} &=& 100.\\ {\rm Dpc} &=& 100.\\ {\rm Dpch} &=& 70.\\ {\rm Tcl} &=& 20.0\\ {\rm Tcl} &=& 20.0\\ {\rm Tch} &=& 15.0\\ {\rm Tch} &=& 2.0\\ {\rm Tal} &=& 0.\\ {\rm Tal} &=& 0.\\ {\rm Tah} &=& 0.\\ {\rm Tah} &=& 0.\\ {\rm Ecl} &=& 1.0\\ {\rm Ech} &=& 1.0\\ {\rm Ech} &=& 0.38\\ {\rm Frh} &=& 0.38\\ {\rm Frh} &=& 0.38\\ {\rm To} &=& 288.5\\ {\rm ITE} &=& 0\\ {\rm Opt} &=& 0.\\ {\rm Opt} &=& 0.3\\ {\rm ITE} &=& 0\\ {\rm Opt} &=& 0.3\\ {\rm Prh} &=& 0.3\\ {\rm Prh} &=& 0.3\\ {\rm Prh} &=& 0.3\\ {\rm Opt} &=& 0.3\\ {\rm $	Gam	-	6.5
$\begin{array}{rcl} Ac &=& 0.54\\ CO2 &=& 0.54\\ H2O &=& 2.17\\ O3 &=& .31\\ CH4 &=& .0009\\ rh0 &=& 0.77\\ Pc1 &=& 0.76\\ Pcm &=& 0.76\\ Pcm &=& 0.56\\ Pch &=& 0.20\\ Dpc &=& 100\\ Dpc &=& 100\\ Dpc &=& 100\\ Dpc &=& 50\\ Dpch &=& 70\\ Tc1 &=& 20.0\\ Tc1 &=& 20.0\\ Tc1 &=& 20.0\\ Tc1 &=& 20.0\\ Tc1 &=& 0.3\\ Tc1 &=& 0\\ Ta1 &=& 0\\ Ta1 &=& 0\\ Ta1 &=& 0\\ Ta1 &=& 0\\ Tch &=& 1.0\\ Cc1 &=& 1.0\\ Ec1 &=& 1.0\\ Ec1 &=& 1.0\\ Ec1 &=& 1.0\\ Ec1 &=& 0.3\\ Fr1 &=& 0.38\\ Frn &=& 0.28\\ Frh &=& 0.38\\ To &=& 288.5\\ Tcn &=& 0.3\\ Tcn &=& $	Pt	-	200.
$\begin{array}{rcl} \text{CO2} &=& 0.54\\ \text{H2O} &=& 2.17\\ \text{O3} &=& .31\\ \text{CH4} &=& .0009\\ \text{rh0} &=& 0.77\\ \text{Pc1} &=& 0.76\\ \text{Pcm} &=& 0.56\\ \text{Pcm} &=& 0.20\\ \text{Dpc} &=& 100\\ \text{Dpc} &=& 100\\ \text{Dpc} &=& 100\\ \text{Dpc} &=& 50\\ \text{Dpch} &=& 70\\ \text{Tc1} &=& 20\\ \text{O}\\ \text{Tc1} &=& 20\\ \text{Tcm} &=& 15\\ 0\\ \text{Tch} &=& 2.0\\ \text{Ta1} &=& 0\\ \text{Tam} &=& 0\\ \text{Camposite} &=& 1.0\\ \text{Ech} &=& 0.38\\ \text{Frn} &=& 0.28\\ \text{Frh} &=& 0.38\\ \text{To} &=& 288.5\\ \text{TTE} &=& 0\\ \end{array}$	Ac		0.54
$\begin{array}{rcl} H20 =& 2.17\\ 03 =& .31\\ CH4 =& .0009\\ rh0 =& 0.77\\ Pc1 =& 0.76\\ Pcm =& 0.56\\ Pch =& 0.20\\ Dpc =& 100\\ Dpc =& 100\\ Dpc =& 100\\ Dpcm =& 50\\ Dpch =& 70\\ Tc1 =& 20.0\\ Tc1 =& 20.0\\ Tc1 =& 2.0\\ Tc1 =& 2.0\\ Tch =& 2.0\\ Tch =& 0.3\\ Tch =& 0\\ Tah =& 0\\ Ec1 =& 1.0\\ Ech =& 0.3\\ Fr1 =& 0.34\\ Frm =& 0.28\\ Frh =& 0.38\\ To =& 288.5\\ Tch =& 288.5\\ Tch =& 0.38\\ Tch =& 288.5\\ Tch =& 0.38\\ Tch =& 0.38\\$	CO2	-	0.54
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	H20		2.17
CH4 =.0009 $rh0 =$ 0.77 $Pcl =$ 0.76 $Pcm =$ 0.56 $Pch =$ 0.20 $Dpc =$ 100. $Dpcm =$ 50. $Dpch =$ 70. $Tcl =$ 20.0 $Tch =$ 2.0 $Tch =$ 2.0 $Tah =$ 0. $Tah =$ 0. $Ecl =$ 1.0 $Ech =$ 0.3 $Frl =$ 0.34 $Frm =$ 0.28 $Frh =$ 0.38 $To =$ 288.5 $TTE =$ 0.34	03		.31
rh0 = $0.77$ $Pcl =$ $0.76$ $Pcm =$ $0.56$ $Pch =$ $0.20$ $Dpc =$ $100.$ $Dpc =$ $100.$ $Dpch =$ $70.$ $Tcl =$ $20.0$ $Tch =$ $2.0$ $Tch =$ $2.0$ $Tah =$ $0.$ $Tah =$ $0.$ $Ecl =$ $1.0$ $Ech =$ $0.34$ $Frn =$ $0.28$ $Frh =$ $0.38$ $To =$ $288.5$	CH4		.0009
Pc1 = $0.76$ $Pcm =$ $0.56$ $Pch =$ $0.20$ $Dpc =$ $100.$ $Dpcm =$ $50.$ $Dpch =$ $70.$ $Tcl =$ $20.0$ $Tch =$ $2.0$ $Tah =$ $0.$ $Tah =$ $0.$ $Ecl =$ $1.0$ $Ech =$ $0.34$ $Frn =$ $0.28$ $Frh =$ $0.38$ $To =$ $288.5$	∙rh0		0.77
Pcm = $0.56$ $Pch =$ $0.20$ $Dpc =$ $100.$ $Dpcm =$ $50.$ $Dpch =$ $70.$ $Tcl =$ $20.0$ $Tch =$ $2.0$ $Tch =$ $2.0$ $Tah =$ $0.$ $Tah =$ $0.$ $Ecl =$ $1.0$ $Ech =$ $0.34$ $Frm =$ $0.28$ $Frh =$ $0.38$ $To =$ $288.5$	Pcl	=	0.76
Pch = $0.20$ Dpc = $100.$ Dpcm = $50.$ Dpch = $70.$ Tcl = $20.0$ Tcm = $15.0$ Tch = $2.0$ Tal = $0.$ Tal = $0.$ Tam = $0.$ Tah = $0.$ Ecl = $1.0$ Ech = $0.34$ Frn = $0.28$ Frh = $0.38$ To = $288.5$ TTE = $0.32$	Pcm	=	0.56
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	Pch	-	0.20
$\begin{array}{rcl} \text{Dpcm} = & 50.\\ \text{Dpch} = & 70.\\ \text{Tcl} = & 20.0\\ \text{Tcm} = & 15.0\\ \text{Tch} = & 2.0\\ \text{Tal} = & 0.\\ \text{Tal} = & 0.\\ \text{Tam} = & 0.\\ \text{Tam} = & 0.\\ \text{Tam} = & 0.\\ \text{Tam} = & 0.\\ \text{Campa } & 1.0\\ \text{Ecl} = & 1.0\\ \text{Ech} = & 1.0\\ \text{Ech} = & 0.3\\ \text{Frl} = & 0.34\\ \text{Frm} = & 0.28\\ \text{Frh} = & 0.38\\ \text{To} = & 288.5\\ \text{Tam} = & 0.\\ \text{Campa } & 0.\\ C$	Dpc	-	100.
$ \begin{array}{rcl} Dpch = & 70. \\ Tcl = & 20.0 \\ Tcm = & 15.0 \\ Tch = & 2.0 \\ Tal = & 0. \\ Tan = & 0. \\ Tam = & 0. \\ Tam = & 0. \\ Tam = & 0. \\ Ccl = & 1.0 \\ Ecl = & 1.0 \\ Ech = & 0.3 \\ Frl = & 0.38 \\ Frn = & 0.28 \\ Frm = & 0.38 \\ To = & 288.5 \\ TTE = & 0 \\ \end{array} $	Dpcn	n	50.
$\begin{array}{rcrcrc} Tcl &=& 20.0\\ Tcm &=& 15.0\\ Tch &=& 2.0\\ Tal &=& 0.\\ Tam &=& 0.\\ Tam &=& 0.\\ Tam &=& 0.\\ Tah &=& 0.\\ Ecl &=& 1.0\\ Ech &=& 1.0\\ Ech &=& 0.3\\ Frl &=& 0.34\\ Frm &=& 0.28\\ Frh &=& 0.38\\ To &=& 288.5\\ TTF &=& 0.\\ TTF &=& 0.\\ TTT &=& 0.\\ Ttr &=& 0.\\ Tt$	Dpch	1 =	70.
$\begin{array}{rcrcrc} Tcm = & 15.0\\ Tch = & 2.0\\ Tal = & 0.\\ Tam = & 0.\\ Tam = & 0.\\ Tah = & 0.\\ Ecl = & 1.0\\ Ecm = & 1.0\\ Ecm = & 1.0\\ Ech = & 0.3\\ Frl = & 0.34\\ Frm = & 0.28\\ Frh = & 0.38\\ To = & 288.5\\ ITE = & 0\end{array}$	Tcl	-	20.0
$\begin{array}{rcl} Tch = & 2.0\\ Tal = & 0.\\ Tam = & 0.\\ Tah = & 0.\\ Ecl = & 1.0\\ Ecm = & 1.0\\ Ech = & 0.3\\ Frl = & 0.34\\ Frm = & 0.28\\ Frh = & 0.38\\ To = & 288.5\\ ITE = & 0\end{array}$	Tcm		15.0
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	$\mathbf{Tch}$	-	2.0
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	Tal	-	0.
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	Tam	÷===	0.
Ec1 = 1.0 Ecm = 1.0 Ech = 0.3 Fr1 = 0.34 Frm = 0.28 Frh = 0.38 To = 288.5 ITE = 0	Tah	<b></b>	0.
$\begin{array}{rcl} Ecm = & 1.0\\ Ech = & 0.3\\ Frl = & 0.34\\ Frm = & 0.28\\ Frh = & 0.38\\ To = & 288.5\\ ITE = & 0 \end{array}$	Ecl	_	1.0
$\begin{array}{rcl} Ech = & 0.3 \\ Fr1 = & 0.34 \\ Frm = & 0.28 \\ Frh = & 0.38 \\ To = & 288.5 \\ ITE = & 0 \end{array}$	Ecm	_	1.0
Fr1 = 0.34 Frm = 0.28 Frh = 0.38 To = 288.5 LTE = 0	Ech	-	0.3
Frm = 0.28 Frh = 0.38 To = 288.5 ITE = 0	Frl	202	0.34
Frh = 0.38 To = 288.5	Frm	_	0.28
$T_0 = 288.5$	Frh	-	0.38
	То	_	288.5
	ITF		0

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# Output for a 3 cloud model, with ITF = 0 Local conditions Filename = SRCOUTB

# SIMPLE CLIMATE MODEL

S0	AMU	DAY	$\mathbf{RG}$	FLUX	сс	WH2O	WCO2	WO3	WCH4	LAP	РТ	EG	RH0
1367.	.50	.50	.13	341.75	.54	2.17	.54	.31	.00090	6.5	200.	1.00	.77

# CLOUD PROPERTIES

	LOW	MIDDLE	HIGH
PC	.76	. 56	.20
DPC	100.	50.0	70.0
CLT	20.0	15.0	2.0
CLA	0.0	0.0	0.0
EC	1.00	1.00	. 30
FR	.34	.28	. 38

ITERATION = 1 TG = 288.5 BETA = 2.9203

LOW CLOUD

TG	TC		TT	TAUG	TAUC
288.5	273.83	212.	42	1.837	1.1492
FIP = .47265					
FIP = .33076					
FIP =7892	1				
FIP = .6362					
FIP =6768	7				
FIP =5058	9				
TBS	TBC				
.63195	.72445				
IROUT CLEAR	= 248.21	IROUT	CLOUD	= 230.92	
IROUT TOTAL	- 238.87				
EBS	EBC				
.21079	.07100				
FBS	FBC		CC	FIRG	
82.792	27.885		. 54	53.142	
MIDDLE CLOUD					
TG	TC		TT	TAUG	TAUC
288.5	258.37	212.	.42	1.837	.71758
FIP = .47265					
FIP = .20649					
FIP =7892	1				
FIP = .6236					
FIP =7700	19				
FIP =2934	-8				
TBS	TBC				
.63195	.80892				
IROUT CLEAR	= 248.21	IROUT	CLOUD	= 204.38	
IROUT TOTAL	= 224.54				
EBS	EBC				
. 21079	.10332				
FBS	FBC		CC	FIRG	
82.792	40.582		, 54	59.998	
HIGH CLOUD					
$\mathbf{TG}$	TC		TT	TAUG	TAUC

288.5 FIP = .47265 212.42 212.42 1.837 .24864 FIP = 6.47093E-2FIP = -.78921FIP = .51168FIP = -.78639FIP = -6.47185E-2TBS TBC .63195 5.7263 IROUT CLEAR = 248.21 IROUT CLOUD = 198.29 IROUT TOTAL = 221.25EBS EBC .21079 .1911 FBS FIRG FBC CC 82.792 75.059 .54 78.616 Κ SUM1 SUM2 DT ΤG ER 228.16 1 227.89 ~288.5 288.5 .001226 NET IR OUT = 228.16 NET SOL IN = 227.89 NET SOLG DOWN = 162.21 NET IRG UP = 64.742GROUND EMISSION = 392.77 ITF=0: SURFACE TEMPERATURE SPECIFIED

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#### VERTICAL STRUCTURE

Z	Р	Т	TAP	TACB	TAC	TAUD
0.0	1000.00	288,50	0.000	1.7800	1.7912	1.8370
.1	990.00	287.95	.3635	1.7473	1.7587	1.8046
. 2	980.00	287.39	.4984	1.7150	1.7265	1.7725
. 3	970.00	286.83	. 5995	1,6830	1.6946	1.7408
. 3	960.00	286.27	.6829	1.6513	1,6630	1.7094
.4	950.00	285.70	.7548	1,6199	1.6318	1.6783
. 5	940.00	285.12	.8185	1.5889	1.6009	1.6475
. 6	930,00	284.54	.8757	1.5582	1.5703	1.6170
. 7	920.00	283,96	.9278	1.5278	1,5400	1.5869
. 8	910.00	283.37	.9756	1.4977	1.5100	<b>1</b> .5571
.9	900.00	282.78	1.0199	1.4679	1.4804	1.5277
1.0	890.00	282.18	1.0610	1.4385	1.4511	1.4985
1.1	880.00	281.57	1.0993	1.4093	1.4221	1.4697
1.2	870.00	280.96	1.1352	1.3805	1.3935	1.4412
1.3	860.00	280.34	1.1690	1.3520	1.3651	1.4131
1.4	850.00	279.72	1.2007	1.3238	1.3371	1,3852
1.4	840.00	279.09	1.2307	1.2959	1.3094	1.3577
1.5	830.00	278.45	1.2590	1.2683	1.2820	1.3305
1.6	820.00	277.81	1.2858	1.2411	1.2549	1,3036
1.7	810.00	277.16	1.3112	1.2142	1.2282	1.2771
1.8	800.00	276.51	1.3353	1.1875	1,2017	1.2509
1.9	790.00	275.85	1.3582	1.1612	1,1756	1.2250
2.0	780.00	275.18	1.3800	1.1352	1.1498	1.1994
2.2	770.00	274.51	1.4007	1.1095	1.1243	1.1741
2.3	760.00	273.83	1.4203	1.0841	<b>1</b> .0991	1.1492
2.4	750.00	273.14	1.4391	1.0590	1.0743	1.1246
2.5	740.00	272.44	1.4569	1.0342	1.0497	1,1003
2.6	730.00	271.74	1.4739	1.0097	1.0255	1.0763
2.7	720.00	271.02	1,4901	.9855	1.0016	1.0527

2.8	710.00	270.30	1.5056	.9617	.9780	1.0294
29	700 00	269 58	1 5203	.9381	9547	1.0064
3 0	690.00	268 84	1 5344	9148	9317	9837
2.0	680.00	268.09	1 5478	8018	9090	9613
J. I 2 2	670.00	200.02	1.5475	9601	. 2020	
3.3	670.00	207.34	1.5005	.0091	.0000	. 9 . 9 . 9 . 2
3.4	660.00	266.08	1.5/2/	.8467	.8646	.91/5
3.5	650.00	265.80	1.5843	.8246	.8428	.8961
3.6	640.00	265.02	1.5954	.8028	.8214	.8/50
3.7	630.00	264.23	1.6060	.7812	.8002	.8542
3.9	620.00	263.42	1.6161	.7600	.7794	.8337
4.0	610.00	262.61	1.6257	.7390	.7588	.8136
4.1	600.00	261.79	1.6348	.7183	.7386	.7938
4.2	590.00	260.95	1.6436	.6979	.7186	.7742
4.4	580.00	260.10	1,6519	.6778	. 6990	.7550
4.5	570.00	259.24	1.6598	.6579	.6796	.7362
4.6	560.00	258.37	1.6674	.6383	.6605	.7176
48	550 00	257.49	1.6746	6189	6418	6993
49	540 00	256 59	1 6814	5998	6233	6814
5 0	530.00	255 68	1 6879	5810	6051	6637
5.0	520.00	255,00	1 6941	5624	5871	6464
53	510.00	254.70	1 7000	5440	5695	6294
5.5	500.00	252.02	1.7000	5258	5521	6107
5.5	500.00	252.00	1.7050	5070	5251	5062
J.0 E 0	490.00	201.09	1,7110	. 5079		. 5902
5.8	480.00	250.91	1.7160	.4902	, 5162	.5602
5.9	470.00	249,90	1.7209	.4/2/	.5017	. 2644
6.L	460.00	248.88	1.7254	.4554	.4854	.5489
6.3	450.00	247.85	1.7298	.4383	.4694	,5337
6.4	440.00	246.79	1.7339	.4213	.4536	.5188
6.6	430.00	245.71	1.7379	.4045	.4381	.5043
6.8	420.00	244.61	1.7416	.3877	.4228	.4900
6.9	410.00	243.49	1.7451	.3711	.4077	.4760
7.1	400.00	242.35	1.7485	.3546	. 3929	.4623
7.3	390.00	241.19	1.7517	. 3380	.3783	.4490
7.5	380.00	240.00	1.7547	.3215	.3639	.4359
7.6	370.00	238.79	1.7576	. 3049	.3496	.4231
7.8	360.00	237.55	1.7604	.2881	.3356	.4106
8.0	350.00	236.28	1.7630	.2712	.3217	. 3984
8.2	340.00	234.98	1.7655	.2538	. 3080	.3865
8.4	330,00	233.65	1.7678	.2359	.2943	. 3749
8.6	320.00	232.28	1.7701	.2172	.2808	.3636
8.9	310.00	230.88	1.7722	.1972	.2672	.3525
9 1	300 00	229.45	1.7743	.1753	.2537	.3417
<u><u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u></u></u>	290.00	227 97	1 7763	1500	2402	3312
05	280.00	226.46	1 7782	1171	2265	3210
9.5	230.00	220.40	1 7800		2125	3111
9.0 10 0	270.00	224.90	1 7017	11/0	1092	3014
10.0	260.00	223.23	1 703/	1/30	1932	2020
10.3	250,00	221.03	1,7034	.1430	.1052	. 2920
10.6	240.00	219.91	1.7051	. 1039	.10/5	. 2020
1U.8	230.00	Z10,14	1 7000	.1/90	.1490	. 2/39
11.1	220.00	216.30	1.7882	. 1923	.1296	.2652
11.4	210,00	214.40	1./89/	.2032	.1030	.2568
11.7	200.00	212.42	1./912	.2125	.0000	.2486
12.0	190.00	212.42	1.7927	. 2207	.1012	.2407
12.4	180.00	212.42	1.7942	.2279	.1245	.2329
12.7	170.00	212.42	1.7957	.2344	.1404	.2254
13.1	160.00	212.42	1.7972	. 2402	.1528	.2180
13.5	150.00	212,42	1.7987	.2456	.1629	.2108

13.9	140.00	212.42	1.8002	.2505	.1715	.2038
14.4	130.00	212.42	1.8018	.2550	.1790	.1969
14.9	120.00	212.42	1.8035	.2593	.1857	.1901
15.4	110.00	212.42	1.8052	.2634	.1917	.1833
16.0	100.00	212.42	1.8070	.2673	.1972	.1766
16.7	90.00	212.42	1.8090	.2712	.2024	.1698
17.4	80.00	212.42	1.8111	.2750	.2073	.1628
18.2	70.00	212.42	1.8133	.2788	.2121	.1556
19.2	60.00	212.42	1.8158	.2828	.2169	.1480
20.3	50.00	212.42	1.8186	.2869	.2218	.1398
21.7	40.00	212.42	1.8216	.2912	.2268	.1305
23.5	30.00	212.42	1.8251	.2959	.2320	.1196
26.0	20.00	212.42	1.8289	.3009	.2375	.1059
30.3	10.00	212.42	1.8331	.3062	.2433	.0863
100.	0.00	212.42	1.8370	.3111	.2486	0.0000

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Alligator Rivers Region Research Institute Research Report 1983-84 Alligator Rivers Region Research Institute Annual Research Summary 1984-85 Alligator Rivers Region Research Institute Annual Research Summary 1985-86 Alligator Rivers Region Research Institute Annual Research Summary 1986-87 Alligator Rivers Region Research Institute Annual Research Summary 1987-88

Research Reports (RR) and Technical Memoranda (TM)

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RR 1	The macroinvertebrates of Magela Creek, Northern Territory. April 1982 (pb, mf - 46 pp.)	R. Marchant
RR 2	Water quality characteristics of eight billabongs in the Magela Creek catchment. December 1982 (pb, mf ~ 60 pp.)	B.T. Hart & R.J. McGregor
RR 3	A limnological survey of the Alligator Rivers Region. I. Diatoms (Bacillariophyceae) of the Region. August 1983 (pb, mf - 160 pp.)	D.P. Thomas
	*A limnological survey of the Alligator Rivers Region. II. Freshwater algae, exclusive of diatoms. 1986 (pb, mf - 176 pp.)	H.U. Ling & P.A. Tyler
RR 4	*Ecological studies on the freshwater fishes of the Alligator Rivers Region, Northern Territory. Volume I. Outline of the study, summary, conclusions and recommendations. 1986 (pb, mf - 63 pp.)	K.A. Bishop, S.A. Allen, D.A. Pollard & M.G. Cook
	Ecological studies on the freshwater fishes of the Alligator Rivers Region, Northern Territory. Volume II. (in press)	K.A. Bishop, S.A. Allen, D.A. Pollard & M.G. Cook
	Ecological studies on the freshwater fishes of the Alligator Rivers Region, Northern Territory. Volume III. (in press)	K.A. Bishop, S.A. Allen, D.A. Pollard & M.G. Cook
RR 5	Macrophyte vegetation of the Magela Creek flood plain, Alligator Rivers Region, Northern Territory. March 1989 (pb - 41 pp.)	C.M. Finlayson, B.J. Bailey & I.D. Cowie
TM 1	Transport of trace metals in the Magela Creek system, Northern Territory. I. Concentrations and loads of iron, manganese, cadmium, copper, lead and zinc during flood periods in the 1978-1979 Wet season. December 1981 (pb, mf - 27 pp.)	B.T. Hart, S.H.R. Davies & P.A. Thomas
TM 2	Transport of trace metals in the Magela Creek system, Northern Territory. II. Trace metals in the Magela Creek billabongs at the end of the 1978 Dry season. December 1981 (pb, mf - 23 pp.)	S.H.R. Davies & B.T. Hart
ТМ 3	Transport of trace metals in the Magela Creek system, Northern Territory. III. Billabong sediments. December 1981 (pb, mf - 24 pp.)	P.A. Thomas, S.H.R. Davies & B.T. Hart
TM 4	The foraging behaviour of herons and egrets on the Magela Creek flood plain, Northern Territory. March 1982 (pb, mf - 20 pp.)	H.R. Recher & R.T. Holmes
TM 5	Flocculation of retention pond water. May 1982 (pb, mf ~ 8 pp.)	B.T. Hart & R.J. McGregor
TM 6	Dietary pathways through lizards of the Alligator Rivers Region Northern Territory. July 1984 (pb, mf - 15 pp.)	C.D. James, S.R. Morton, R.W. Braithwaite & J.C. Wombey
<b>TM</b> 7	Capacity of waters in the Magela Creek system, Northern Territory, to complex copper and cadmium. August 1984 (pb, mf - 42 pp.)	B.T. Hart & S.H.R. Davies
TM 8	Acute toxicity of copper and zinc to three fish species from the Alligator Rivers Region. August 1984 (pb, mf - 31 pp.)	L. Baker & D. Walden

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ТМ 9	Textural characteristics and heavy metal concentrations in billabong sediments from the Magela Creek system, northern Australia. October 1984 (pb, mf - 39 pp.)	P.A. Thomas & B.T. Hart	
TM 10	Oxidation of manganese(II) in Island Billabong water. October 1984 (pb, mf ~ 11 pp.)	B.T. Hart & M.J. Jones	
<b>TM</b> 11	<i>In situ</i> experiments to determine the uptake of copper by the aquatic macrophyte <i>Najas tenuifolia</i> R.Br. December 1984 (pb, mf - 13 pp.)	B.T. Hart, M.J. Jones & P. Breen	
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TM 13	Fate, of heavy metals in the Magela Creek system, northern Australia. I. Experiments with plastic enclosures placed in Island Billabong during the 1980 Dry Season: heavy metals. May 1985 (pb, mf - 46 pp.)	B.T. Hart, M.J. Jones & P. Bek	
TM 14	Fate of heavy metals in the Magela Creek system, northern Australia. II. Experiments with plastic enclosures placed in Island Billabong during the 1980 Dry season: limnology and phytoplankton.	B.T. Hart, M.J. Jones, P. Bek & J. Kessell	
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